

A degenerating Cahn-Hilliard system coupled with complete damage processes

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Abstract

In this work, we analytically investigate a degenerating PDE system for phase separation and complete damage processes considered on a nonsmooth time-dependent domain with mixed boundary conditions. The evolution of the system is described by a *degenerating* Cahn-Hilliard equation for the concentration, a doubly nonlinear differential inclusion for the damage variable and a *degenerating* quasi-static balance equation for the displacement field. All these equations are highly nonlinearly coupled. Because of the doubly degenerating character of the system, the doubly nonlinear differential inclusion and the nonsmooth domain, the structure of the model is very complex from an analytical point of view.

A novel approach is introduced for proving existence of weak solutions for such degenerating coupled system. To this end, we first establish a suitable notion of weak solutions, which consists of weak formulations of the diffusion and the momentum balance equation, a variational inequality for the damage process and a total energy inequality. To show existence of weak solutions, several new ideas come into play. Various results on shrinking sets and its corresponding local Sobolev spaces are used. It turns out that, for instance, on open sets which shrink in time a quite satisfying analysis in Sobolev spaces is possible. The presented analysis can handle highly nonsmooth regions where complete damage takes place. To mention only one difficulty, infinitely many completely damaged regions which are not connected with the Dirichlet boundary may occur in arbitrary small time intervals.

Key Words: Cahn-Hilliard system, phase separation, complete damage, elliptic-parabolic degenerating systems, linear elasticity, weak solution, doubly nonlinear differential inclusions, existence results, rate-dependent systems.

AMS Subject Classifications: 35K85, 35K55, 49J40, 49S05, 35J50, 74A45, 74G25, 34A12, 82B26, 82C26, 35K92, 35K65, 35K35;

1 Problem description

Phase separation and damage processes occur in many fields, including material sciences, biology and chemical reactions. In particular, for the manufacturing and lifetime prediction of micro-electronic devices it is of great importance to understand the mechanisms and the interplay between phase-separation and damage processes in solder alloys. As soon as elastic alloys are quenched sufficiently, spinodal decomposition leads to a fine-grained structure of different chemical mixtures on a short time-scale (see [DM01] for numerical simulations and experimental observations). The long-term evolution is determined by a chemical diffusion process which tends to minimize the bulk and the

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surface energy of the chemical substances. J.W. Cahn and J.E. Hilliard developed a phenomenological model for the kinetics of phase-separation in a thermodynamically consistent framework known as the Cahn-Hilliard equation [CH58], for which an extensive mathematical literature exists. An overview of modeling and analytical aspects of the Cahn-Hilliard equation can be found in [Ell89]. The recent literature is mainly focused on coupled systems. For instance, physical observations and numerical simulations reveal that mechanical stresses influence the developing of shapes of the chemical phases. A coupling between Cahn-Hilliard systems and elastic deformations have been analytically studied in [Gar00, BCD⁺02, CMP00, Gar05a, Gar05b, BP05, PZ08]. For numerical results and simulations we refer to [Wei01, Mer05, BB99, GRW01, BM10]. Phase separation of the chemical components may also lead to critical stresses at phase boundaries due to swelling which result in cracks and formation of voids and are of particular interest to understand the aging process in solder materials, cf. [HCW91, USG07, GUE⁺07, FK09]. A fully coupled system consisting of the Cahn-Hilliard equation, an elliptic equation for the displacement field and a differential inclusion for the damage variable has been recently investigated in [HK11, HK13]. However, in [HK11, HK13] and in the most mathematical damage literature [BS04, Gia05, MR06, MT10, KRZ11], it is usually assumed that damage cannot completely disintegrate the material (i.e. *incomplete damage*). Dropping this assumption gives rise to many mathematical challenges. Therefore, global-in-time existence results for complete damage models are rare. Modeling and existence of weak solutions for *purely* mechanical complete damage systems with quasi-static force balances are studied in [BMR09, Mie11, HK12] and with visco-elasticity in [MRZ10, RR12].

The main goal in the present work is to prove existence of weak solutions of a system *coupling* a differential inclusion describing damage processes with an elastic Cahn-Hilliard system in a small strain setting as in [HK11] but allowing for *complete damage* and *degenerating mobilities* with respect to the damage variable (see Theorem 4.5). Note that the mobility is still assumed to be independent of the concentration variable (see [EG96] for an analysis of Cahn-Hilliard equations with degenerating and concentration dependent mobilities). The elasticity is considered to be linear and the system is assumed to be in quasi-static mechanical equilibrium since diffusion processes take place on a much slower time scale. The main modeling idea in [HK12] has been to formulate such degenerating system on a time-dependent domain which consists of the not completely damaged regions that are connected to the Dirichlet boundary. In this context, the concept of maximal admissible subsets is introduced to specify the domain of interest.

In the following, we fix a bounded \mathcal{C}^2 -domain $\Omega \subseteq \mathbb{R}^n$ and a Dirichlet boundary part $D \subseteq \partial\Omega$ with $\mathcal{H}^{n-1}(D) > 0$.

A relatively open subset G of $\overline{\Omega}$ is called *admissible* with respect to an \mathcal{H}^{n-1} -measurable part D of the boundary $\partial\Omega$ if every path-connected component P_G of G satisfies $\mathcal{H}^{n-1}(P_G \cap D) > 0$, where \mathcal{H}^{n-1} denotes the $(n - 1)$ -dimensional Hausdorff measure. The *maximal admissible subset* of G is denoted by $\mathfrak{A}_D(G)$. With the notion of maximal admissible subsets, we can formulate our evolutionary system with a time-dependent domain.

Degenerating PDE system on a time-dependent domain.

Find functions

$$c \in \mathcal{C}^2(F; \mathbb{R}), \quad u \in \mathcal{C}_x^2(F; \mathbb{R}^n), \quad z \in \mathcal{C}^2(\overline{\Omega_T}; \mathbb{R}), \quad \mu \in \mathcal{C}_x^2(F; \mathbb{R})$$

with the time-dependent domain $F = \{(x, t) \in \overline{\Omega_T} \mid x \in F(t)\}$ and $F(t) = \mathfrak{A}_D(\{z(t) > 0\})$ such that

the PDE system

$$\begin{aligned}
0 &= \operatorname{div}(W_{,e}(c, \epsilon(u), z)), \\
c_t &= \operatorname{div}(m(z) \nabla \mu), \\
\mu &= -\Delta c + \Psi_{,c}(c) + W_{,c}(c, \epsilon(u), z), \\
z_t + \xi - \operatorname{div}(|\nabla z|^{p-2} \nabla z) + W_{,z}(c, \epsilon(u), z) + f'(z) &= 0, \\
\xi &\in \partial I_{(-\infty, 0]}(z_t)
\end{aligned}$$

is satisfied pointwise in $\operatorname{int}(F)$ with the initial-boundary conditions

$$\begin{aligned}
c(0) = c^0, \quad z(0) = z^0 & \quad \text{in } F(0), \\
u(t) = b(t) & \quad \text{on } \Gamma_1(t) := F(t) \cap D, \\
W_{,e}(c(t), \epsilon(u(t)), z(t)) \cdot \nu = 0 & \quad \text{on } \Gamma_2(t) := F(t) \cap (\partial\Omega \setminus D), \\
z(t) = 0 & \quad \text{on } \Gamma_3(t) := \partial F(t) \setminus F(t), \\
\nabla z(t) \cdot \nu = 0 & \quad \text{on } \Gamma_1(t) \cup \Gamma_2(t), \\
\nabla c(t) \cdot \nu = 0 & \quad \text{on } \Gamma_1(t) \cup \Gamma_2(t), \\
m(z(t)) \nabla \mu(t) \cdot \nu = 0 & \quad \text{on } \Gamma_1(t) \cup \Gamma_2(t).
\end{aligned}$$

Here, $\mathcal{C}_x^2(F; \mathbb{R}^N)$ denotes the space of two times continuously differentiable functions with respect to the spatial variable x on the set F .

The solution of the PDE system can physically be interpreted as follows: c denotes the chemical concentration difference of a two-component alloy, $\epsilon(u) := \frac{1}{2}(\nabla u + (\nabla u)^T)$ the linearized strain tensor of the displacement u , z the damage profile describing the degree of damage (i.e. $z = 1$ undamaged and $z = 0$ completely damaged material point) and μ the chemical potential. Moreover, W denotes the elastic energy density, Ψ the chemical energy density, f a damage dependent potential, m the mobility depending on the damage variable z and b the time-dependent Dirichlet boundary data for D . *Capillarity effects* are modeled by concentration gradients in the free energy which provokes the chemical potential to depend on the Laplacian Δc (see [Gur89]).

We assume the following product structure for the elastic energy density:

$$W(c, e, z) = g(z) \varphi(c, e) \tag{1}$$

with a non-negative monotonically increasing function $g \in \mathcal{C}^1([0, 1]; \mathbb{R}^+)$ such that the complete damage condition $g(0) = 0$ is fulfilled. The second function $\varphi \in \mathcal{C}^1(\mathbb{R} \times \mathbb{R}_{\text{sym}}^{n \times n}; \mathbb{R}^+)$ should have the following polynomial form

$$\varphi(c, e) = \varphi^1 e : e + \varphi^2(c) : e + \varphi^3(c) \tag{2}$$

for coefficients $\varphi^1 \in \mathcal{L}(\mathbb{R}_{\text{sym}}^{n \times n})$ with φ^1 positive definite, $\varphi^2 \in \mathcal{C}^1(\mathbb{R}; \mathbb{R}_{\text{sym}}^{n \times n})$ and $\varphi^3 \in \mathcal{C}^1(\mathbb{R})$. Note that homogeneous elastic energy densities of the type

$$W(c, e, z) = \frac{1}{2} z \mathbb{C} (e - e^*(c)) : (e - e^*(c))$$

with eigenstrain e^* depending on the concentration and stiffness tensor \mathbb{C} are covered in this approach. We remark that this approach allows to incorporate swelling phenomena of the chemical

species which, in turn, can lead to critical stresses for the initiation of damage. The mobility $m \in \mathcal{C}([0, 1]; \mathbb{R}^+)$, on the other hand, should satisfy the following condition for degeneracy, i.e.

$$m(z) = 0 \text{ if and only if } z = 0. \quad (3)$$

Plan of the paper

A weak formulation of the degenerating system is given in Section 2 while the proof of existence is carried out in Section 3 and Section 4. In the first step of the proof, a degenerated limit of the corresponding incomplete damage system coupled with elastic Cahn-Hilliard equations is performed (see Section 3). Due to the additional coupling the passage to the limit becomes quite involved and a *conical Poincaré inequality* is used to control the chemical potential in a local sense.

In the next step, we take material exclusions into account (see Section 4). They occur when not completely damaged components become disconnected to the Dirichlet boundary. The exclusions lead to jumps in the energy inequality. By a concatenation property and by Zorn's lemma, we are able to prove the main theorem in this paper: a global-in-time existence result (see Theorem 4.5). In addition, several Γ -results and various properties on shrinking sets and their corresponding local Sobolev spaces are used.

2 Assumptions and notion of weak solutions

As before, $\Omega \subseteq \mathbb{R}^n$ denotes a bounded \mathcal{C}^2 -domain and $D \subseteq \partial\Omega$ the Dirichlet boundary with $\mathcal{H}^{n-1}(D) > 0$. Furthermore, $T > 0$ denotes the maximal time of interest and $p > n$ a fixed exponent.

The weak formulation of the PDE system presented in this work will be based on an energetic approach and uses the associated free energy. Let $G \subseteq \overline{\Omega}$ be a relatively open subset. Then, the free energy \mathcal{E} on G is given by

$$\mathcal{E}_G(c, e, z) := \int_G \left(\frac{1}{p} |\nabla z|^p + \frac{1}{2} |\nabla c|^2 + \Psi(c) + W(c, e, z) + f(z) \right) dx$$

for $c \in H^1(G)$, $e \in L^2(G; \mathbb{R}_{\text{sym}}^{n \times n})$ and $z \in W^{1,p}(G)$.

We suppose the structural assumptions (1) and (2). Here, $\Psi \in C^1(\mathbb{R})$ denotes the chemical energy density, $f \in C^1([0, 1])$ a damage dependent potential and $g \in C^1([0, 1]; \mathbb{R}^+)$ with $g(0) = 0$ a function modeling the influence of the damage on the material stiffness. We assume the following growth conditions:

$$|\varphi^2(c)|, |\varphi_{,c}^2(c)| \leq C(1 + |c|), \quad (4a)$$

$$|\varphi^3(c)|, |\varphi_{,c}^3(c)| \leq C(1 + |c|^2), \quad (4b)$$

$$|\Psi_{,c}(c)| \leq C(1 + |c|^{2^*/2}), \quad (4c)$$

$$\eta \leq g'(z). \quad (4d)$$

The constants $\eta, C > 0$ are independent of c and z , and 2^* denotes the Sobolev critical exponent. In the case $n = 2$, $\Psi_{,c}$ has to satisfy an r -growth condition for a fixed arbitrary $r > 0$ whereas we have no restrictions on $\Psi_{,c}$ in the one-dimensional case.

The functions $m \in \mathcal{C}([0, 1]; \mathbb{R}^+)$ will denote the mobility of the Cahn-Hilliard diffusion process and is assumed to satisfy (3).

Before we state the precise notion of weak solutions, we introduce some notation.

The set $\{f > 0\}$ for a given function $f \in W^{1,p}(\Omega)$ has to be read as $\{x \in \overline{\Omega} \mid f(x) > 0\}$ by employing the embedding $W^{1,p}(\Omega) \hookrightarrow \mathcal{C}(\overline{\Omega})$ (because of $p > n$). The symbol $H_S^1(G)$ denotes the subspace of $H^1(G)$ with zero trace on S , where $G \subset \mathbb{R}^n$ is open and S is an \mathcal{H}^{n-1} -measurable subset of $\partial\Omega$.

Let $I \subseteq \mathbb{R}$ be an open interval. The subspace $SBV(I; X) \subseteq BV(I; X)$ of special functions of bounded variation is defined as the space of functions $f \in BV(I; X)$ where the decomposition

$$df = f' \mathcal{L}^1 + (f^+ - f^-) \mathcal{H}^0 \llcorner J_f$$

for an $f' \in L^1(I; X)$ exists. The function f' is called the absolutely continuous part of the differential measure and we also write $\partial_t^a f$ and J_f denotes the jump set. If, additionally, $\partial_t^a f \in L^p(I; X)$, $p \geq 1$, we write $f \in SBV^p(I; X)$.

We say that a relatively open subset $F \subseteq \overline{\Omega_T}$ is shrinking iff $F(t) \subseteq F(s)$ for all $0 \leq s \leq t \leq T$, where $F(t)$ denotes the t -cut of F . For the analysis of our proposed system, we introduce local Sobolev functions. The space-time local Sobolev space $L_t^2 H_{x,\text{loc}}^1(F; \mathbb{R}^N)$ for a shrinking set F is given by

$$L_t^2 H_{x,\text{loc}}^q(F; \mathbb{R}^N) := \left\{ v : F \rightarrow \mathbb{R}^N \mid \forall t \in (0, T], \forall U \subset\subset F(t) \text{ open} : v|_{U \times (0,t)} \in L^2(0, t; H^q(U; \mathbb{R}^N)) \right\}.$$

In the following, a weak formulation of the system above combining the ideas in [HK11] and [HK12] is given.

Definition 2.1 (Weak solution of the coupled PDE system) *A quadruple (c, u, z, μ) is called a weak solution with the initial-boundary data (c^0, z^0, b) if*

(i) *Spaces:*

$$\begin{aligned} c &\in L^\infty(0, T; H^1(\Omega)) \cap H^1(0, T; (H^1(\Omega))^*), & c(0) &= c^0, \\ u &\in L_t^2 H_{x,\text{loc}}^1(F; \mathbb{R}^n), & u &= b \text{ on } D_T \cap F, \\ z &\in L^\infty(0, T; W^{1,p}(\Omega)) \cap SBV^2(0, T; L^2(\Omega)), & z(0) &= z^0, \quad z^+(t) = z^-(t) \mathbf{1}_{F(t)} \\ \mu &\in L_t^2 H_{x,\text{loc}}^1(F) \end{aligned}$$

with $e := \epsilon(u) \in L^2(F; \mathbb{R}_{\text{sym}}^{n \times n})$ where $F := \mathfrak{A}_D(\{z^- > 0\}) \subseteq \overline{\Omega_T}$ is a shrinking set ($z^+(t) := \lim_{s \downarrow t} z(s)$ and $z^-(t) := \lim_{s \uparrow t} z(s)$ in $L^2(\Omega)$).

(ii) *Quasi-static mechanical equilibrium:*

$$0 = \int_{F(t)} W_{,e}(c(t), e(t), z(t)) : \epsilon(\zeta) \, dx \quad (5)$$

for a.e. $t \in (0, T)$ and for all $\zeta \in H_D^1(\Omega; \mathbb{R}^n)$.

(iii) *Diffusion:*

$$\int_{\Omega_T} \partial_t \zeta (c - c^0) \, dx \, dt = \int_F m(z) \nabla \mu \cdot \nabla \zeta \, dx \, dt \quad (6)$$

for all $\zeta \in L^2(0, T; H^1(\Omega)) \cap H^1(0, T; L^2(\Omega))$ with $\zeta(T) = 0$ and

$$\int_F \mu \zeta \, dx \, dt = \int_F \left(\nabla c \cdot \nabla \zeta + \Psi_{,c}(c) \zeta + W_{,c}(c, e, z) \zeta \right) dx \, dt \quad (7)$$

for all $\zeta \in L^2(0, T; H^1(\Omega))$ with $\text{supp}(\zeta) \subseteq F$.

(iv) *Damage variational inequality:*

$$\begin{aligned} 0 &\leq \int_{F(t)} \left(|\nabla z(t)|^{p-2} \nabla z(t) \cdot \nabla \zeta + (W_{,z}(c(t), e(t), z(t)) + f'(z(t)) + \partial_t^a z(t)) \zeta \right) dx \\ 0 &\leq z(t), \\ 0 &\geq \partial_t^a z(t) \end{aligned} \quad (8)$$

for a.e. $t \in (0, T)$ and for all $\zeta \in W^{1,p}(\Omega)$ with $\zeta \leq 0$.

(v) *Energy inequality:*

$$\mathcal{E}(t) + \mathcal{J}(0, t) + \mathcal{D}(0, t) \leq \mathcal{E}^+(0) + \mathcal{W}_{\text{ext}}(0, t), \quad (9)$$

where the terms in the energy inequality are as follows:

- *Energy:* $\mathcal{E}(t) := \mathcal{E}_{F(t)}(c(t), e(t), z(t)),$
- *Energy jump term:* $\mathcal{J}(0, t) := \sum_{s \in J_z \cap (0, t]} (\mathcal{E}^-(s) - \mathcal{E}^+(s)),$
 $\mathcal{E}^-(t) := \lim_{s \rightarrow t^-} \left(\text{ess inf}_{\tau \in (s, t)} \mathcal{E}(\tau) \right),$
 $\mathcal{E}^+(t) := \mathfrak{e}_t^+$
with $0 \leq \mathfrak{e}_t^+ \leq \inf_{\zeta \in H_{D \cap F(t)}^1(F(t); \mathbb{R}^n)} \mathcal{E}_{F(t)}(c(t), \epsilon(b(t) + \zeta), z(t)),$
- *Dissipated energy:* $\mathcal{D}(0, t) := \int_0^t \int_{F(s)} \left(m(z) |\nabla \mu|^2 + |\partial_t^a z|^2 \right) dx \, dt,$
- *External work:* $\mathcal{W}_{\text{ext}}(0, t) := \int_0^t \int_{F(s)} W_{,e}(c, \epsilon(u), z) : \epsilon(\partial_t b) \, dx \, dt.$

Remark 2.2 Under additional regularity assumptions, a weak solution satisfies

$$\begin{aligned} \lim_{\tau \rightarrow s^-} \mathcal{E}_{F(\tau)}(c(\tau), e(\tau), z(\tau)) &= \lim_{\tau \rightarrow s^-} \text{ess inf}_{\vartheta \in (\tau, s)} \mathcal{E}_{F(\vartheta)}(c(\vartheta), e(\vartheta), z(\vartheta)), \\ \lim_{\tau \rightarrow s^+} \mathcal{E}_{F(\tau)}(c(\tau), e(\tau), z(\tau)) &= \mathfrak{e}_s^+ \end{aligned}$$

and reduces to the pointwise classical formulation presented in Section 1 (cf. [HK12, Theorem 3.7]).

As mentioned before, the main aim of this paper is to prove global-in-time existence of weak solutions. Due to the analytical quite involved structure, we introduce the notion of weak solutions with a given fineness constant $\eta > 0$.

Definition 2.3 (Weak solution with fineness η) A tuple (c, e, u, z, μ) together with a shrinking set $F \subseteq \overline{\Omega_T}$ is called a weak solution with fineness $\eta > 0$ if

$$\begin{aligned} c &\in L^\infty(0, T; H^1(\Omega)) \cap H^1(0, T; (H^1(\Omega))^*), & c(0) &= c^0, \\ u &\in L_t^2 H_{x, \text{loc}}^1(\mathfrak{A}_D(F); \mathbb{R}^n), & u &= b \text{ on } D_T \cap \mathfrak{A}_D(F), \\ z &\in L^\infty(0, T; W^{1,p}(\Omega)) \cap SBV^2(0, T; L^2(\Omega)), & z(0) &= z^0, \quad z^+(t) = z^-(t) \mathbf{1}_{F(t)} \\ \mu &\in L_t^2 H_{x, \text{loc}}^1(F) \end{aligned}$$

such that $e = \epsilon(u)$ in $\mathfrak{A}_D(F)$ and the properties (ii)-(v) of Definition 2.1 and

$$\begin{aligned} \forall t \in [0, T] : & \quad \mathfrak{A}_D(\{z^-(t) > 0\}) \subseteq F(t) \text{ and } \mathcal{L}^n(F(t) \setminus \mathfrak{A}_D(\{z^-(t) > 0\})) < \eta, \\ \forall t \in [0, T] \setminus \bigcup_{t \in C_{z^*}} [t, t + \eta) & \quad \mathfrak{A}_D(\{z^-(t) > 0\}) = F(t). \end{aligned}$$

are satisfied.

Here, C_{z^*} denotes the set of cluster points from the right of the jump set J_{z^*} of the function $z^* \in SBV^2(0, T; L^2(\Omega))$ given by $z^*(t) := z(t) \mathbf{1}_{\mathfrak{A}_D(\{z^-(t) > 0\})}$, i.e., z^* is the damage profile of z restricted to the time-dependent domain $\mathfrak{A}_D(\{z^- > 0\})$ ($\mathbf{1}_A : X \rightarrow \{0, 1\}$ denotes the characteristic function of a set $A \subseteq X$).

Remark 2.4 (i) If a weak solution (c, e, u, z, μ) on F with fineness η according to Definition 2.3 satisfies $C_{z^*} = \emptyset$ then (c, u, z, μ) is a weak solution according to Definition 2.1.

(ii) If the initial value for the damage profile contains no completely damaged parts on $\overline{\Omega}$, we are also able to establish maximal local-in-time existence of weak solutions. More precisely, there exist a maximal value $\hat{T} > 0$ with $\hat{T} \leq T$ and functions c, u, z and μ defined on the time interval $[0, \hat{T}]$ such that (c, u, z, μ) is a weak solution according to Definition 2.1. Therefore, if $\hat{T} < T$, (c, u, z, μ) cannot be extended to a weak solution on $[0, \hat{T} + \varepsilon]$.

The remaining part of this paper is devoted to establish a global existence result in the sense of Definition 2.3 (see Theorem 4.5).

3 Degenerate limit of the regularized system

In this section, we will start with a corresponding incomplete damage model coupled to an elastic Cahn-Hilliard system and then perform a limit procedure. For each $\varepsilon > 0$, we define the regularized free energy \mathcal{E}_ε as

$$\mathcal{E}_\varepsilon(c, e, z) := \int_{\Omega} \left(\frac{1}{p} |\nabla z|^p + \frac{1}{2} |\nabla c|^2 + \Psi(c) + W^\varepsilon(c, e, z) + f(z) \right) dx$$

for functions $c \in H^1(\Omega)$, $e \in L^2(\Omega; \mathbb{R}_{\text{sym}}^{n \times n})$, $z \in W^{1,p}(\Omega)$. The regularized elastic energy density and mobility are given by

$$\begin{aligned} W^\varepsilon(c, e, z) &:= (g(z) + \varepsilon) \varphi(c, e), \\ m^\varepsilon(z) &:= m(z) + \varepsilon. \end{aligned}$$

A modification of the proof of Theorem 4.6 in [HK11] yields the following result.

Theorem 3.1 (ε -regularized coupled PDE problem) *Let $\varepsilon > 0$. For given initial-boundary data $c_\varepsilon^0 \in H^1(\Omega)$, $z_\varepsilon^0 \in W^{1,p}(\Omega)$ and $b_\varepsilon \in W^{1,1}(0, T; W^{1,\infty}(\Omega; \mathbb{R}^n))$ there exists a quadruple $q_\varepsilon = (c_\varepsilon, u_\varepsilon, z_\varepsilon, \mu_\varepsilon)$ such that*

(i) *Spaces:*

$$\begin{aligned} c_\varepsilon &\in L^\infty(0, T; H^1(\Omega)) \cap H^1(0, T; (H^1(\Omega))^*), & c_\varepsilon(0) &= c_\varepsilon^0, \\ u_\varepsilon &\in L^\infty(0, T; H^1(\Omega; \mathbb{R}^n)), & u_\varepsilon &= b_\varepsilon \text{ on } D_T, \\ z_\varepsilon &\in L^\infty(0, T; W^{1,p}(\Omega)) \cap H^1(0, T; L^2(\Omega)), & z_\varepsilon(0) &= z_\varepsilon^0, \\ \mu_\varepsilon &\in L^2(0, T; H^1(\Omega)) \end{aligned}$$

(ii) *Quasi-static mechanical equilibrium:*

$$\int_{\Omega} W_{,e}^\varepsilon(c_\varepsilon(t), \epsilon(u_\varepsilon(t)), z_\varepsilon(t)) : \epsilon(\zeta) \, dx = 0 \quad (10)$$

for a.e. $t \in (0, T)$ and for all $\zeta \in H_D^1(\Omega; \mathbb{R}^n)$.

(iii) *Diffusion:*

$$\int_{\Omega_T} (c_\varepsilon - c_\varepsilon^0) \partial_t \zeta \, dx \, dt = \int_{\Omega_T} m^\varepsilon(z_\varepsilon) \nabla \mu_\varepsilon \cdot \nabla \zeta \, dx \, dt \quad (11)$$

for all $\zeta \in L^2(0, T; H^1(\Omega))$ with $\partial_t \zeta \in L^2(\Omega_T)$ and $\zeta(T) = 0$ and

$$\int_{\Omega} \mu_\varepsilon(t) \zeta \, dx = \int_{\Omega} \left(\nabla c_\varepsilon(t) \cdot \nabla \zeta + \Psi_{,c}(c_\varepsilon(t)) \zeta + W_{,c}^\varepsilon(c_\varepsilon(t), \epsilon(u_\varepsilon(t)), z_\varepsilon(t)) \zeta \right) dx \quad (12)$$

for a.e. $t \in (0, T)$ and for all $\zeta \in H^1(\Omega)$.

(iv) *Damage variational inequality:*

$$0 \leq \int_{\Omega} \left(|\nabla z_\varepsilon(t)|^{p-2} \nabla z_\varepsilon(t) \cdot \nabla \zeta + (W_{,z}^\varepsilon(c_\varepsilon(t), \epsilon(u_\varepsilon(t)), z_\varepsilon(t)) + f'(z_\varepsilon(t)) + \partial_t z_\varepsilon(t) + r_\varepsilon(t)) \zeta \right) dx \quad (13)$$

$$0 \leq z_\varepsilon(t),$$

$$0 \geq \partial_t z_\varepsilon(t)$$

for a.e. $t \in (0, T)$ and for all $\zeta \in W^{1,p}(\Omega)$ with $\zeta \leq 0$, where $r_\varepsilon \in L^1(\Omega_T)$ is given by

$$r_\varepsilon = -\chi_\varepsilon (W_{,z}^\varepsilon(c_\varepsilon, \epsilon(u_\varepsilon), z_\varepsilon) + f'(z))^{+} \quad (14)$$

with $\chi_\varepsilon \in L^\infty(\Omega)$ fulfilling $\chi_\varepsilon = 0$ on $\{z_\varepsilon > 0\}$ and $0 \leq \chi_\varepsilon \leq 1$ on $\{z_\varepsilon = 0\}$.

(v) *Energy inequality:*

$$\begin{aligned} \mathcal{E}_\varepsilon(c_\varepsilon(t), \epsilon(u_\varepsilon(t)), z_\varepsilon(t)) + \int_{\Omega_t} \left(m^\varepsilon(z_\varepsilon) |\nabla \mu_\varepsilon|^2 + |\partial_t z_\varepsilon|^2 \right) dx \, ds \\ \leq \mathcal{E}_\varepsilon(c_\varepsilon^0, \epsilon(u_\varepsilon^0), z_\varepsilon^0) + \int_{\Omega_t} W_{,e}(c_\varepsilon, \epsilon(u_\varepsilon), z_\varepsilon) : \epsilon(\partial_t b_\varepsilon) \, dx \, ds \end{aligned} \quad (15)$$

holds for a.e. $t \in (0, T)$ where u_ε^0 minimizes $\mathcal{E}_\varepsilon(c^0, \epsilon(\cdot), z_\varepsilon^0)$ in $H^1(\Omega; \mathbb{R}^n)$ with Dirichlet data $b_\varepsilon^0 := b_\varepsilon(0)$ on D .

Proof. The existence theorem presented in [HK11] can be adapted to our situation by considering the viscous semi-implicit time-discretized system (in a classical notation; we omit the ε -dependence in the notation for the discrete solution at the moment):

$$\begin{aligned} 0 &= \operatorname{div}(W_{,e}^\varepsilon(c^k, \epsilon(u^k), z^k)) + \delta \operatorname{div}(|\nabla u^k|^2 \nabla u^k), \\ \frac{c^k - c^{k-1}}{\tau} &= \operatorname{div}(m^\varepsilon(z^{k-1}) \nabla \mu^k), \\ \mu^k &= -\Delta c^k + \Psi_{,c}(c^k) + W_{,c}^\varepsilon(c^k, \epsilon(u^k), z^k) + \delta \frac{c^k - c^{k-1}}{\tau}, \\ \frac{z^k - z^{k-1}}{\tau} + \xi + \zeta &= \operatorname{div}(|\nabla z^k|^{p-2} \nabla z^k) + W_{,z}^\varepsilon(c^k, \epsilon(u^k), z^k) + f'(z^k), \end{aligned}$$

with the sub-gradients $\xi \in \partial I_{(-\infty, 0]}((z^k - z^{k-1})/\tau)$, $\zeta \in \partial I_{[0, \infty)}(z^k)$ and the discretization fineness $\tau = T/M$ for $M \in \mathbb{N}$. The discrete equations can be obtained recursively starting from (c^0, u^0, z^0) with $u^0 := \arg \min_{u \in H^1(\Omega; \mathbb{R}^n), u|_D = b^0|_D} \mathcal{E}_\varepsilon(c^0, u, z^0)$ by considering the Euler-Lagrange equations of the functional

$$\begin{aligned} \mathbb{E}^k(c, u, z) &:= \mathcal{E}_\varepsilon(c, u, z) + \int_\Omega \frac{\delta}{4} |\nabla u|^4 dx \\ &\quad + \frac{\tau}{2} \left(\left\| \frac{z - z^{k-1}}{\tau} \right\|_{L^2(\Omega)}^2 + \left\| \frac{c - c^{k-1}}{\tau} \right\|_{X(z^{k-1})}^2 + \delta \left\| \frac{c - c^{k-1}}{\tau} \right\|_{L^2(\Omega)}^2 \right) \end{aligned}$$

defined on the subspace of $H^1(\Omega) \times W^{1,4}(\Omega; \mathbb{R}^n) \times W^{1,p}(\Omega)$ with the conditions $u|_D = b(k\tau)|_D$, $\int_\Omega (c - c^0) dx = 0$ and $0 \leq z \leq z^{k-1}$ a.e. in Ω . The weighted scalar product $\langle \cdot, \cdot \rangle_{X(z^{k-1})}$ is given by

$$\langle u, v \rangle_{X(z^{k-1})} := \left\langle m^\varepsilon(z^{k-1}) \nabla A^{-1} u, \nabla A^{-1} v \right\rangle_{L^2(\Omega)}$$

with the operator $A : V_0 \rightarrow \tilde{V}_0$, $Au := \langle m^\varepsilon(z^{k-1}) \nabla u, \nabla \cdot \rangle_{L^2(\Omega)}$ and the spaces

$$\begin{aligned} V_0 &:= \left\{ \zeta \in H^1(\Omega) \mid \int_\Omega \zeta dx = 0 \right\}, \\ \tilde{V}_0 &:= \left\{ \zeta \in (H^1(\Omega))^* \mid \langle \zeta, \mathbf{1} \rangle_{(H^1)^* \times H^1} = 0 \right\}. \end{aligned}$$

After passing the discretization fineness to 0, i.e. $\tau \rightarrow 0^+$, we obtain the corresponding regularized equations and inequalities for (10)-(15). A further passage $\delta \rightarrow 0^+$ yields a weak solution as required. \square

In the following, we always consider a given countable sequence $\varepsilon_k \rightarrow 0^+$ as $k \rightarrow \infty$ and omit the subscript k by simply writing $\varepsilon \rightarrow 0^+$.

The existence proof is based on Γ -limit techniques of the *reduced* energy functionals of \mathcal{E}_ε and \mathcal{F}_ε with $\mathcal{F}_\varepsilon(c, e, z) := \int_\Omega W^\varepsilon(c, e, z) dx$ in order to gain a suitable energy estimate in the limit $\varepsilon \rightarrow 0^+$. The reduced energy functionals \mathfrak{E}_ε and \mathfrak{F}_ε are defined by

$$\mathfrak{E}_\varepsilon(c, \xi, z) := \begin{cases} \min_{\zeta \in H_D^1(\Omega; \mathbb{R}^n)} \mathcal{E}_\varepsilon(c, \epsilon(\xi + \zeta), z) & \text{if } 0 \leq z \leq 1, \\ \infty & \text{else,} \end{cases}$$

$$\mathfrak{F}_\varepsilon(c, \xi, z) := \begin{cases} \min_{\zeta \in H_D^1(\Omega; \mathbb{R}^n)} \mathcal{F}_\varepsilon(c, \epsilon(\xi + \zeta), z) & \text{if } 0 \leq z \leq 1, \\ \infty & \text{else.} \end{cases}$$

The Γ -limits of \mathfrak{E}_ε and \mathfrak{F}_ε as $\varepsilon \rightarrow 0^+$ exist in the topological space $H_w^1(\Omega) \times W^{1,\infty}(\Omega; \mathbb{R}^n) \times W_w^{1,p}(\Omega)$ and are denoted by \mathfrak{E} and \mathfrak{F} , respectively. Here, $H_w^1(\Omega)$ denotes the space $H^1(\Omega)$ with its weak topology. The limit functional \mathfrak{F} is needed as an auxiliary construction in the following because it already captures the essential properties of \mathfrak{E} . In the next section, we are going to prove some properties of the Γ -limit \mathfrak{E} which are used in the global-in-time existence proof.

Let $(c_\varepsilon^0, b_\varepsilon^0, z_\varepsilon^0) \rightarrow (c^0, b^0, z^0)$ as $\varepsilon \rightarrow 0^+$ be a recovery sequence for $\mathfrak{E}_\varepsilon \xrightarrow{\Gamma} \mathfrak{E}$. In particular, $c_\varepsilon^0 \rightharpoonup c^0$ in $H^1(\Omega)$, $b_\varepsilon^0 \rightarrow b^0$ in $W^{1,\infty}(\Omega; \mathbb{R}^n)$ and $z_\varepsilon^0 \rightharpoonup z^0$ in $W^{1,p}(\Omega)$. Furthermore, we set $b_\varepsilon := b - b^0 + b_\varepsilon^0$. For each $\varepsilon > 0$, we obtain a weak solution $(c_\varepsilon, u_\varepsilon, z_\varepsilon, \mu_\varepsilon)$ for $(c_\varepsilon^0, z_\varepsilon^0, b_\varepsilon)$ according to Theorem 3.1.

By means of the energy estimate (15), we deduce the following a-priori estimates for our proposed model:

Lemma 3.2 *There exists a constant $C > 0$ independent of ε such that*

- (i) $\sup_{t \in [0, T]} \|c_\varepsilon(t)\|_{H^1(\Omega)} \leq C,$
- (ii) $\|\widehat{e}_\varepsilon\|_{L^2(\Omega_T; \mathbb{R}^{n \times n})} \leq C$ with $\widehat{e}_\varepsilon := e_\varepsilon \mathbf{1}_{\{z_\varepsilon > 0\}},$
- (iii) $\sup_{t \in [0, T]} \|z_\varepsilon(t)\|_{W^{1,p}(\Omega)} \leq C,$
- (iv) $\|\partial_t z_\varepsilon\|_{L^2(\Omega_T)} \leq C,$
- (v) $\|W^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon)\|_{L^\infty(0, T; L^1(\Omega))} \leq C$ with $e_\varepsilon := \epsilon(u_\varepsilon),$
- (vi) $\|m^\varepsilon(z_\varepsilon)^{1/2} \nabla \mu_\varepsilon\|_{L^2(\Omega_T; \mathbb{R}^n)} \leq C,$
- (vii) $\|\partial_t c_\varepsilon\|_{L^2(0, T; (H^1(\Omega))^*)} \leq \|m^\varepsilon(z_\varepsilon) \nabla \mu_\varepsilon\|_{L^2(\Omega_T; \mathbb{R}^n)} \leq C.$

Proof. Applying Gronwall's lemma to the energy estimate (15) and noticing the boundedness of $\mathcal{E}_\varepsilon(c_\varepsilon^0, \epsilon(u_\varepsilon^0), z_\varepsilon^0)$ with respect to $\varepsilon \in (0, 1)$ show (iv), (vi) and

$$\mathcal{E}_\varepsilon(c_\varepsilon(t), e_\varepsilon(t), z_\varepsilon(t)) \leq C \tag{16}$$

for a.e. $t \in (0, T)$ and all $\varepsilon \in (0, 1)$ and in particular (v). Item (vii) follows from (vi) and the diffusion equation (11). Taking the constraint $\int_\Omega c_\varepsilon(t) dx \equiv \text{const}$ and the restriction $0 \leq z_\varepsilon \leq 1$ into account, property (16) gives rise to

$$\begin{aligned} \|c_\varepsilon\|_{L^\infty(0, T; H^1(\Omega))} &\leq C, \\ \|z_\varepsilon\|_{L^\infty(0, T; W^{1,p}(\Omega))} &\leq C. \end{aligned}$$

Together with the control of the time-derivatives (iv) and (vii), we even obtain boundedness of

$$\begin{aligned} \|c_\varepsilon(t)\|_{H^1(\Omega)} &\leq C, \\ \|z_\varepsilon(t)\|_{W^{1,p}(\Omega)} &\leq C \end{aligned}$$

for every $t \in [0, T]$ and $\varepsilon \in (0, 1)$. Hence, (i) and (iii) are proven.

It remains to show (ii). By using the estimates

$$\begin{aligned} 0 &\leq \chi_\varepsilon \leq 1 && \text{a.e. in } \Omega_T, \\ 0 &\leq W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) && \text{a.e. in } \Omega_T, \end{aligned}$$

and the definition of r_ε in (14), we obtain (here $(\cdot)^+ := \max\{\cdot, 0\}$ and $(\cdot)^- := \min\{\cdot, 0\}$)

$$\begin{aligned} &\int_{\{z_\varepsilon > 0\}} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon)) \, dx \, dt \\ &= \int_{\Omega_T} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon)) \, dx \, dt \\ &\quad - \int_{\{z_\varepsilon = 0\}} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon))^+ \, dx \, dt - \int_{\{z_\varepsilon = 0\}} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon))^- \, dx \, dt \\ &\leq \int_{\Omega_T} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon)) \, dx \, dt \\ &\quad - \int_{\{z_\varepsilon = 0\}} \chi_\varepsilon (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon))^+ \, dx \, dt - \int_{\{z_\varepsilon = 0\}} (f'(z_\varepsilon))^- \, dx \, dt \\ &= \int_{\Omega_T} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon) + r_\varepsilon) \, dx \, dt - \int_{\{z_\varepsilon = 0\}} (f'(z_\varepsilon))^- \, dx \, dt. \end{aligned} \tag{17}$$

To proceed, we test inequality (13) with $\zeta \equiv -1$ and integrate from $t = 0$ to $t = T$:

$$\int_{\Omega_T} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) + f'(z_\varepsilon) + r_\varepsilon) \, dx \, dt \leq - \int_{\Omega_T} \partial_t z_\varepsilon \, dx \, dt. \tag{18}$$

Combining (17) and (18), we end up with

$$\int_{\{z_\varepsilon > 0\}} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon)) \, dx \, dt \leq - \int_{\Omega_T} \partial_t z_\varepsilon \, dx \, dt - \int_{\{z_\varepsilon = 0\}} (f'(z_\varepsilon))^- \, dx \, dt.$$

The right hand side is bounded because of the boundedness properties (iii) and (iv). But we also have ($\tilde{C} > 0$ and $C > 0$ are constants)

$$\begin{aligned} \int_{\{z_\varepsilon > 0\}} (W_{,z}^\varepsilon(c_\varepsilon, e_\varepsilon, c_\varepsilon) + f'(z_\varepsilon)) \, dx \, dt &\geq \int_{\{z_\varepsilon > 0\}} (\tilde{C}|e_\varepsilon|^2 - C|c_\varepsilon|^2 + f'(z_\varepsilon)) \, dx \, dt \\ &\geq \tilde{C} \int_{\Omega_T} |\widehat{e}_\varepsilon|^2 \, dx \, dt + \int_{\{z_\varepsilon > 0\}} (-C|c_\varepsilon|^2 + f'(z_\varepsilon)) \, dx \, dt. \end{aligned}$$

Since the left hand side is bounded and $\int_{\{z_\varepsilon > 0\}} (-C|c_\varepsilon|^2 + f'(z_\varepsilon)) \, dx \, dt$ is also bounded, we obtain (ii). \square

We proceed by deriving convergence properties for c_ε , \widehat{e}_ε , z_ε as $\varepsilon \rightarrow 0$.

Lemma 3.3 *There exists functions*

$$c \in L^\infty(0, T; H^1(\Omega)) \cap H^1(0, T; (H^1(\Omega))^*) \text{ with } c(0) = c^0,$$

$$\widehat{e} \in L^2(\Omega_T; \mathbb{R}^{n \times n}),$$

$$z \in L^\infty(0, T; W^{1,p}(\Omega)) \cap H^1(0, T; L^2(\Omega)) \text{ with } z(0) = z^0, \quad z \leq 0, \quad \partial_t z \leq 0$$

and a subsequence (we omit the index) such that for $\varepsilon \rightarrow 0^+$

$$\begin{aligned} (a) \quad & c_\varepsilon \rightharpoonup c \text{ in } H^1(0, T; (H^1(\Omega))^*), \\ & c_\varepsilon \rightarrow c \text{ in } L^\kappa(0, T; L^r(\Omega)) \text{ for all } \kappa \geq 1, \quad 1 \leq r < 2^*, \\ & c_\varepsilon(t) \rightarrow c(t) \text{ in } H^1(\Omega) \text{ for all } t, \\ & c_\varepsilon \rightarrow c \text{ a.e. in } \Omega_T, \end{aligned}$$

$$\begin{aligned} (b) \quad & z_\varepsilon \rightharpoonup z \text{ in } H^1(0, T; L^2(\Omega)), \\ & z_\varepsilon \rightarrow z \text{ in } L^\kappa(0, T; W^{1,p}(\Omega)) \text{ for all } \kappa \geq 1, \\ & z_\varepsilon(t) \rightharpoonup z(t) \text{ in } W^{1,p}(\Omega) \text{ for all } t, \\ & z_\varepsilon \rightarrow z \text{ in } \mathcal{C}(\overline{\Omega_T}), \end{aligned}$$

$$(c) \quad b_\varepsilon \rightarrow b \text{ in } W^{1,1}(0, T; W^{1,\infty}(\Omega; \mathbb{R}^n)),$$

$$\begin{aligned} (d) \quad & \widehat{e}_\varepsilon \rightharpoonup \widehat{e} \text{ in } L^2(\Omega_T; \mathbb{R}^{n \times n}), \\ & W_{,e}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) \rightharpoonup W_{,e}(c, \widehat{e}, z) \text{ in } L^2(\{z > 0\}; \mathbb{R}^{n \times n}), \\ & W_{,e}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) \rightarrow 0 \text{ in } L^2(\{z = 0\}; \mathbb{R}^{n \times n}), \\ & W_{,c}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) \rightharpoonup W_{,c}(c, \widehat{e}, z) \text{ in } L^2(\{z > 0\}; \mathbb{R}^{n \times n}), \\ & W_{,c}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) \rightarrow 0 \text{ in } L^2(\{z = 0\}; \mathbb{R}^{n \times n}). \end{aligned}$$

Proof. The a-priori estimates from Lemma 3.2 and classical compactness theorems as well as compactness theorems from Lions and Aubin [Sim86] yield the convergence properties of (a) and

$$z_\varepsilon \xrightarrow{*} z \quad \text{in } L^\infty(0, T; W^{1,p}(\Omega)), \quad (19a)$$

$$z_\varepsilon \rightharpoonup z \quad \text{in } H^1(0, T; L^2(\Omega)), \quad (19b)$$

$$z_\varepsilon \rightarrow z \quad \text{in } \mathcal{C}(\overline{\Omega_T}), \quad (19c)$$

$$\widehat{e}_\varepsilon \rightharpoonup \widehat{e} \quad \text{in } L^2(\Omega_T; \mathbb{R}^{n \times n}), \quad (19d)$$

$$W_{,e}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) \rightharpoonup w_e \quad \text{in } L^2(\Omega_T; \mathbb{R}^{n \times n}) \quad (19e)$$

as $\varepsilon \rightarrow 0^+$ for a subsequence (we omit the subscript) and appropriate functions w_e , \widehat{e} and z .

The proof of the strong convergence of ∇z_ε in $L^p(\Omega_T; \mathbb{R}^n)$ essentially relies on the uniform p -monotonicity property

$$C_{uc}|x - y|^p \leq (|x|^{p-2}x - |y|^{p-2}y) \cdot (x - y) \quad (20)$$

and on an approximation scheme $\{\zeta_\varepsilon\} \subseteq L^p(0, T; W^{1,p}(\Omega))$ with $\zeta_\varepsilon \geq 0$ and

$$\zeta_\varepsilon \rightarrow z \text{ in } L^p(0, T; W^{1,p}(\Omega)) \text{ as } \varepsilon \rightarrow 0^+, \quad (21a)$$

$$0 \leq \zeta_\varepsilon \leq z_\varepsilon \text{ a.e. in } \Omega_T \text{ for all } \varepsilon \in (0, 1). \quad (21b)$$

For details of the construction of ζ_ε , we refer to [HK11]. By using (20), we obtain the estimate:

$$C_{uc} \int_{\Omega_T} |\nabla z_\varepsilon - \nabla z|^p \, dx \, dt \leq \int_{\Omega_T} (|\nabla z_\varepsilon|^{p-2} \nabla z_\varepsilon - |\nabla z|^{p-2} \nabla z) \cdot \nabla (z_\varepsilon - z) \, dx \, dt$$

$$\begin{aligned}
&= \underbrace{\int_{\Omega_T} |\nabla z_\varepsilon|^{p-2} \nabla z_\varepsilon \cdot \nabla (z_\varepsilon - \zeta_\varepsilon) \, dx \, dt}_{A_\varepsilon} \\
&\quad + \underbrace{\int_{\Omega_T} |\nabla z_\varepsilon|^{p-2} \nabla z_\varepsilon \cdot \nabla (\zeta_\varepsilon - z) - |\nabla z|^{p-2} \nabla z \cdot \nabla (z_\varepsilon - z) \, dx \, dt}_{B_\varepsilon}.
\end{aligned}$$

The variational inequality (13) tested with $\zeta(t) = \zeta_\varepsilon(t) - z_\varepsilon(t)$ and integrated from $t = 0$ to $t = T$ yield

$$\begin{aligned}
A_\varepsilon &\leq \int_{\Omega_T} \left(W_{,z}^\varepsilon(c_\varepsilon, \epsilon(u_\varepsilon), z_\varepsilon) + f'(z_\varepsilon) + \partial_t z_\varepsilon \right) (\zeta_\varepsilon - z_\varepsilon) \, dx \, dt \\
&\leq \int_{\Omega_T} g'(z_\varepsilon) \varphi^1 \epsilon(u_\varepsilon) : \epsilon(u_\varepsilon) (\zeta_\varepsilon - z_\varepsilon) \, dx \, dt + \int_{\Omega_T} \varphi^2(c_\varepsilon) : \epsilon(u_\varepsilon) (\zeta_\varepsilon - z_\varepsilon) \, dx \, dt \\
&\quad + \int_{\Omega_T} (\varphi^3(c_\varepsilon) + f'(z_\varepsilon) + \partial_t z_\varepsilon) (\zeta_\varepsilon - z_\varepsilon) \, dx \, dt.
\end{aligned} \tag{22}$$

Here, we have used the fact that $r_\varepsilon = 0$ a.e. in $\{z_\varepsilon > 0\}$ and $\zeta_\varepsilon - z_\varepsilon = 0$ a.e. in $\{z_\varepsilon = 0\}$ by (21b). We treat each term on the right hand side of (22) as follows:

- Due to (4d) and (21b), we obtain

$$g'(z_\varepsilon) \varphi^1 \epsilon(u_\varepsilon) : \epsilon(u_\varepsilon) (\zeta_\varepsilon - z_\varepsilon) \leq 0 \text{ a.e. in } \Omega_T.$$

Therefore, the first integral on the right hand side of (22) is less or equal 0.

- Due to the estimates (4a) and (21b), and the estimates

$$z_\varepsilon \leq Cg(z_\varepsilon),$$

(follows from (4d) and $g(0) \geq 0$) and

$$g(z_\varepsilon) |e_\varepsilon|^2 \leq C(W^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) + |c_\varepsilon|^2)$$

(follows from (2), (4a) and (4a)), the second integral on the right hand side of (22) can be estimated as follows:

$$\begin{aligned}
&\left| \int_{\Omega_T} \varphi^2(c_\varepsilon) : \epsilon(u_\varepsilon) (\zeta_\varepsilon - z_\varepsilon) \, dx \, dt \right| \\
&\leq \|\varphi^2(c_\varepsilon)\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \|\epsilon(u_\varepsilon) \sqrt{|\zeta_\varepsilon - z_\varepsilon|}\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \|\sqrt{|\zeta_\varepsilon - z_\varepsilon|}\|_{L^1(0,T;L^\infty(\Omega))} \\
&\leq C \|c_\varepsilon\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \|\epsilon(u_\varepsilon) \sqrt{|z_\varepsilon| + |\zeta_\varepsilon|}\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \|\sqrt{|\zeta_\varepsilon - z_\varepsilon|}\|_{L^1(0,T;L^\infty(\Omega))} \\
&\leq C \|c_\varepsilon\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \|\epsilon(u_\varepsilon)^2 g(z_\varepsilon)\|_{L^\infty(0,T;L^1(\Omega;\mathbb{R}^{n \times n}))}^{1/2} \|\sqrt{|\zeta_\varepsilon - z_\varepsilon|}\|_{L^1(0,T;L^\infty(\Omega))} \\
&\leq C \|c_\varepsilon\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \left(\|W^\varepsilon(c_\varepsilon, \epsilon(u_\varepsilon), z_\varepsilon)\|_{L^\infty(0,T;L^1(\Omega;\mathbb{R}^{n \times n}))}^{1/2} + \|c_\varepsilon\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \right) \times \\
&\quad \times \|\sqrt{|\zeta_\varepsilon - z_\varepsilon|}\|_{L^1(0,T;L^\infty(\Omega))}.
\end{aligned} \tag{23}$$

By using the a priori estimates in Lemma 3.2 (i) and (v), we obtain

$$\|c_\varepsilon\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \left(\|W^\varepsilon(c_\varepsilon, \epsilon(u_\varepsilon), z_\varepsilon)\|_{L^\infty(0,T;L^1(\Omega;\mathbb{R}^{n \times n}))}^{1/2} + \|c_\varepsilon\|_{L^\infty(0,T;L^2(\Omega;\mathbb{R}^{n \times n}))} \right) \leq C$$

and by using the convergence properties (19c) and (21a) and the compact embedding $W^{1,p}(\Omega) \hookrightarrow \mathcal{C}(\overline{\Omega_T})$, we end up with

$$\|\sqrt{|\zeta_\varepsilon - z_\varepsilon|}\|_{L^1(0,T;L^\infty(\Omega))} \rightarrow 0$$

as $\varepsilon \rightarrow 0^+$. Therefore, the right hand side of (23) converges to 0 as $\varepsilon \rightarrow 0^+$.

- Exploiting the boundedness properties in Lemma 3.2 (i), (iii) and (iv) and the convergence properties (19c) and (21a), the third integral on the right hand side of (22) converges to 0 as $\varepsilon \rightarrow 0^+$.

We have shown $\limsup_{\varepsilon \rightarrow 0^+} A_\varepsilon = 0$. Furthermore, the boundedness of $\{\nabla z_\varepsilon\}$ in $L^p(\Omega_T; \mathbb{R}^n)$ (see Lemma 3.2) and the convergence property (21a) show $B_\varepsilon \rightarrow 0$ as $\varepsilon \rightarrow 0^+$.

We conclude with $\lim_{\varepsilon \rightarrow 0^+} \int_{\Omega_T} |\nabla z_\varepsilon - \nabla z|^p dx dt = 0$. Therefore, (b) is also shown.

Property (c) is an immediate consequence of the definition of b_ε and the convergence property $b_\varepsilon^0 \rightarrow b^0$ in $W^{1,\infty}(\Omega; \mathbb{R}^n)$.

To prove (d), we define $N_\varepsilon := \{z_\varepsilon > 0\} \cap \{z > 0\}$. Consequently, we get

$$W_{,e}^\varepsilon(c_\varepsilon, \widehat{e}_\varepsilon, z_\varepsilon) \mathbb{1}_{N_\varepsilon} = W_{,e}^\varepsilon(c_\varepsilon, e_\varepsilon, z_\varepsilon) \mathbb{1}_{N_\varepsilon} \quad (24)$$

and the convergence

$$\mathbb{1}_{N_\varepsilon} \rightarrow \mathbb{1}_{\{z>0\}} \text{ pointwise in } \Omega_T \text{ and in } L^q(\Omega_T) \text{ for all } q \geq 1 \quad (25)$$

as $\varepsilon \rightarrow 0^+$ by using $z_\varepsilon \rightarrow z$ in $\overline{\Omega_T}$. Calculating the weak $L^1(\Omega_T; \mathbb{R}^{n \times n})$ -limits in (24) for $\varepsilon \rightarrow 0^+$ on both sides by using the already proven convergence properties, we obtain $W_{,e}(c, \widehat{e}, z) \mathbb{1}_{\{z>0\}} = w_e \mathbb{1}_{\{z>0\}}$. The remaining convergence properties in (d) follow from Lemma 3.2 (v). \square

Corollary 3.4 *Let $t \in [0, T]$ and $U \subset \subset \{z(t) = 0\}$ be an open subset. Then $U \subseteq \{z_\varepsilon(s) > 0\}$ for all $s \in [0, t]$ provided that $\varepsilon > 0$ is sufficiently small. More precisely, there exist $0 < \eta$ such that*

$$z_\varepsilon(s) \geq \eta \text{ in } U$$

for all $s \in [0, t]$ and for all $0 < \varepsilon \ll 1$.

Proof. We have $\overline{U} \times [0, T] \subseteq \{z > 0\}$ and $\inf_{(x,t) \in \overline{U} \times [0,t]} z(x, s) > 0$ by $z \in \mathcal{C}(\overline{\Omega_T})$. The claim follows from $z_\varepsilon \rightarrow z$ in $\mathcal{C}(\overline{\Omega_T})$ according to Lemma 3.3 (b). \square

To obtain a-priori estimates for the chemical potentials $\{\mu_\varepsilon\}$ in a local sense, we make use of the so-called *conical* Poincaré inequality for star-shaped domains cited below.

Theorem 3.5 (Conical Poincaré inequality [BK98]) *Suppose that $\Omega \subseteq \mathbb{R}^n$ is a bounded and star-shaped domain, $r \geq 0$ and $1 \leq p < \infty$. Then there exists a constant $C = C(\Omega, p, r) > 0$ such that*

$$\int_{\Omega} |w(x) - w_{\Omega, \delta^r}|^p \delta^r(x) dx \leq C \int_{\Omega} |\nabla w(x)|^p \delta^r(x) dx$$

for all $w \in C^1(\Omega)$, where the δ^r -weight w_{Ω, δ^r} is given by

$$w_{\Omega, \delta^r} := \int_{\Omega} w(x) \delta^r(x) dx, \quad \delta(x) := \text{dist}(x, \partial\Omega).$$

By a density argument, the statement is, of course, also true for all $w \in W^{1,p}(\Omega)$.

Lemma 3.6 (A-priori estimates for μ_ε)

(i) *Interior estimate. For every $t \in [0, T]$ and for every open cube $Q \subset\subset \{z(t) > 0\} \cap \Omega$, there exists a $C > 0$ such that for all $0 < \varepsilon \ll 1$*

$$\|\mu_\varepsilon\|_{L^2(0,t;H^1(Q))} \leq C. \quad (26)$$

(ii) *Estimate at the boundary. For every $t \in [0, T]$ and every $x_0 \in \{z(t) > 0\} \cap \partial\Omega$, there exist a neighborhood U of x_0 and a $C > 0$ such that for all $0 < \varepsilon \ll 1$*

$$\|\mu_\varepsilon\|_{L^2(0,t;H^1(U \cap \Omega))} \leq C. \quad (27)$$

Proof.

(i) Let $t \in [0, T]$ and $Q \subset\subset \{z(t) > 0\} \cap \Omega$ be an open cube. We consider the Lipschitz domain $\tilde{Q} := B_\eta(Q) := \{x \in \mathbb{R}^n \mid \text{dist}(x, Q) < \varepsilon\}$, where $\eta > 0$ is chosen so small such that $\tilde{Q} \subset\subset \{z(t) > 0\} \cap \Omega$. We define the following function

$$\zeta(x) := \begin{cases} \lambda(x) := \text{dist}(x, \partial\tilde{Q}) & \text{if } x \in \tilde{Q}, \\ 0 & \text{else.} \end{cases}$$

ζ is a Lipschitz function on $\bar{\Omega}$ with Lipschitz constant 1. By Rademacher's theorem (see, for instance, [Zie89, Theorem 2.2.1]), ζ is in $W^{1,\infty}(\Omega)$. Now, we can test (12) with ζ . Then, by using the previously proven a-priori estimates, we obtain boundedness of

$$\int_{\tilde{Q}} \mu_\varepsilon(x, s) \lambda(x) dx \leq C \quad (28)$$

with respect to a.e. $s \in (0, T)$ and ε .

Furthermore, there exists an $\eta > 0$ such that $z_\varepsilon(s) \geq \eta$ in \tilde{Q} for all $s \in [0, t]$ and for all $0 < \varepsilon \ll 1$ (see Corollary 3.4). Thus, by assumption (3), $m^\varepsilon(z_\varepsilon(s)) \geq \eta' > 0$ holds in \tilde{Q} for all $s \in [0, t]$ and all $0 < \varepsilon \ll 1$ for a common constant $\eta' > 0$. Consequently, we get by the a-priori estimate for $m^\varepsilon(z_\varepsilon)^{1/2} \nabla \mu_\varepsilon$

$$\|\nabla \mu_\varepsilon\|_{L^2(\tilde{Q} \times [0, t])} \leq C \quad (29)$$

for all ε . Applying Theorem 3.5 (we put in $\Omega = \tilde{Q}$, $r = p = 2$ and $w = \mu_\varepsilon(s)$ for $s \in [0, t]$), integrating from 0 to t in time and using boundedness properties (28) and (29), we obtain boundedness of $\|\mu_\varepsilon \lambda\|_{L^2(\tilde{Q} \times [0, t])}$ and thus boundedness of $\|\mu_\varepsilon\|_{L^2(Q \times [0, t])}$ with respect to $0 < \varepsilon \ll 1$. Together with (29), we get the claim (26).

(ii) Since $\{z(t) = 0\} \subseteq \overline{\Omega}$ is a closed set, we can find a neighborhood $U \subseteq \mathbb{R}^n \setminus \{z(t) = 0\}$ of x_0 . Furthermore, since Ω has a \mathcal{C}^2 -boundary, there exists a \mathcal{C}^2 -diffeomorphism $\pi : (-1, 1)^n \rightarrow U$ with the properties

- $\pi((-1, 1)^{n-1} \times (-1, 0)) \subseteq \Omega$,
- $\pi((-1, 1)^{n-1} \times \{0\}) \subseteq \partial\Omega$,
- $\pi((-1, 1)^{n-1} \times (0, 1)) \subseteq \mathbb{R}^n \setminus \overline{\Omega}$.

Let $\vartheta : (-1, 1)^n \rightarrow (-1, 1)^n$ denote the reflection $x \mapsto (x_1, \dots, x_{n-1}, -x_n)$ and $\mathcal{T} := \pi \circ \vartheta \circ \pi^{-1}$. Furthermore, let $\tilde{\mu}_\varepsilon \in L^2(0, t; H^1(U))$ be defined by

$$\tilde{\mu}_\varepsilon(x, s) := \begin{cases} \mu_\varepsilon(x, s) & \text{if } x \in U \cap \Omega, \\ \mu_\varepsilon(\mathcal{T}(x), s) & \text{if } x \in U \setminus \overline{\Omega}. \end{cases}$$

Let $Q \subset\subset U$ be a non-empty open cube with $x_0 \in Q$. Then, integration by substitution with respect to the transformation \mathcal{T} yields for a.e. $s \in (0, t)$

$$\begin{aligned} \int_Q \tilde{\mu}_\varepsilon(x, s) \lambda(x) \, dx &= \int_{Q \cap \Omega} \mu_\varepsilon(x, s) \lambda(x) \, dx \\ &\quad + \int_{\mathcal{T}(Q \setminus \Omega)} \mu_\varepsilon(x, s) \lambda(\mathcal{T}(x)) |\det(\nabla \mathcal{T}(x))| \, dx, \end{aligned} \quad (30)$$

where the Lipschitz function $\lambda : \mathbb{R}^N \rightarrow \mathbb{R}$ is given by

$$\lambda(x) := \begin{cases} \text{dist}(x, \partial Q) & \text{if } x \in Q, \\ 0 & \text{if } x \in \mathbb{R}^n \setminus Q. \end{cases}$$

We are going to show that both terms on the right hand side of (30) are bounded with respect to ε and a.e. $s \in (0, t)$.

- Testing (12) with the function $\zeta = \lambda$ yields

$$\begin{aligned} \int_{Q \cap \Omega} \mu_\varepsilon(s) \lambda \, dx &= \int_{Q \cap \Omega} (\nabla c_\varepsilon(s) \cdot \nabla \lambda + \Psi_{,c}(c_\varepsilon(s)) \lambda) \, dx \\ &\quad + \int_{Q \cap \Omega} W_{,c}^\varepsilon(c_\varepsilon(s), \epsilon(u_\varepsilon(s)), z_\varepsilon(s)) \lambda \, dx. \end{aligned}$$

By the already know a-priori estimates, every integral term on the right hand side is bounded w.r.t. ε and a.e. $s \in (0, T)$.

- The function

$$\zeta(x) := \begin{cases} (\lambda(\mathcal{T}(x))) |\det(\nabla \mathcal{T}(x))| & \text{if } x \in \mathcal{T}(Q \setminus \Omega), \\ 0 & \text{if } x \in \Omega \setminus \mathcal{T}(Q \setminus \Omega) \end{cases} \quad (31)$$

is a Lipschitz function in Ω because:

- $\lambda \circ \mathcal{T}$ is a Lipschitz function in $U \cap \Omega$ and $\lambda \circ \mathcal{T} = 0$ in $(U \cap \Omega) \setminus \mathcal{T}(Q \setminus \Omega)$. The first property follows from the Lipschitz continuity of λ and of \mathcal{T} (note that \mathcal{T} is a \mathcal{C}^2 -diffeomorphism). The latter property can be seen as follows. *Assume the contrary.* Then, we find an $x \in (U \cap \Omega) \setminus \mathcal{T}(Q \setminus \Omega)$ such that $\lambda(\mathcal{T}(x)) > 0$. By the definition of λ , we get $\mathcal{T}(x) \in Q$. Since $x \in \Omega$, it follows $\mathcal{T}(x) \notin \Omega$ by the construction of \mathcal{T} . Therefore, $\mathcal{T}(x) \in Q \setminus \Omega$. This gives $x = \mathcal{T}(\mathcal{T}(x)) \in \mathcal{T}(Q \setminus \Omega)$ which is a contradiction.
- $|\det(\nabla \mathcal{T})|$ is a Lipschitz function in $U \cap \Omega$ (note that \mathcal{T} is a \mathcal{C}^2 -diffeomorphism).

Testing (12) with ζ from (31) yields

$$\begin{aligned}
& \int_{\mathcal{T}(Q \setminus \Omega)} \mu_\varepsilon(s) (\lambda \circ \mathcal{T}) |\det(\nabla \mathcal{T})| dx \\
&= \int_{\mathcal{T}(Q \setminus \Omega)} \nabla c_\varepsilon(s) \cdot \nabla ((\lambda \circ \mathcal{T}) |\det(\nabla \mathcal{T})|) dx \\
&+ \int_{\mathcal{T}(Q \setminus \Omega)} (\Psi_{,c}(c_\varepsilon(s)) + W_{,c}(c_\varepsilon(s), \epsilon(u_\varepsilon(s)), z_\varepsilon(s))) (\lambda \circ \mathcal{T}) |\det(\nabla \mathcal{T})| dx.
\end{aligned}$$

By the already know a-priori estimates, every integral term on the right hand side is bounded w.r.t. ε and a.e. $s \in (0, t)$.

For $\nabla \tilde{\mu}_\varepsilon(s)$, we also get by integration via substitution:

$$\begin{aligned}
& \int_0^t \int_Q |\nabla \tilde{\mu}_\varepsilon(x, s)|^2 dx ds \\
& \leq \int_0^t \int_{Q \cap \Omega} |\nabla \mu_\varepsilon(x, s)|^2 dx ds + \int_0^t \int_{Q \setminus \Omega} |\nabla \mu_\varepsilon(\mathcal{T}(x), s)|^2 |\nabla \mathcal{T}(x)|^2 dx ds \\
& = \int_0^t \int_{Q \cap \Omega} |\nabla \mu_\varepsilon(x, s)|^2 dx ds \\
& + \int_0^t \int_{\mathcal{T}(Q \setminus \Omega)} |\nabla \mu_\varepsilon(x, s)|^2 |\nabla \mathcal{T}(\mathcal{T}(x))|^2 |\det(\nabla \mathcal{T}(x))| dx ds. \tag{32}
\end{aligned}$$

Since $Q \cap \Omega \subset \subset \{z(t) > 0\}$ and $\mathcal{T}(Q \setminus \Omega) \subset \subset \{z(t) > 0\}$, we deduce $z_\varepsilon(s) \geq \eta$ on $Q \cap \Omega$ and on $\mathcal{T}(Q \setminus \Omega)$ for all $s \in [0, t]$ and for all sufficiently small $0 < \varepsilon$ (see Corollary 3.4). Thus, $\nabla \mu_\varepsilon$ is bounded in $L^2((Q \cap \Omega) \times (0, t); \mathbb{R}^n)$ and in $L^2(\mathcal{T}(Q \setminus \Omega) \times (0, t); \mathbb{R}^n)$ with respect to $0 < \varepsilon \ll 1$ by also using the a-priori estimate for $m^\varepsilon(z_\varepsilon)^{1/2} \nabla \mu_\varepsilon$, the property $m \in \mathcal{C}([0, 1]; \mathbb{R}_+)$ and assumption (3).

Therefore, the left hand side of (32) is also bounded for all $0 < \varepsilon \ll 1$. The Conical Poincaré inequality in Theorem 3.5 yields boundedness of $\tilde{\mu}_\varepsilon \lambda$ in $L^2(Q \times (0, t))$. Finally, we can find a neighborhood $V \subseteq Q$ of x_0 such that $\tilde{\mu}_\varepsilon$ is bounded in $L^2(0, t; H^1(V))$. \square

Due to the a-priori estimates for $\{\mu_\varepsilon\}$ and $\{u_\varepsilon\}$, the limit functions μ and u can only be expected to be in some space-time local Sobolev space $L_t^2 H_{x, \text{loc}}^1$ (see Section 2). In the sequel, it will be necessary to represent the maximal admissible subset of the not completely damaged area, i.e. $\mathfrak{A}_D(\{z > 0\})$, as a union of Lipschitz domains which are connected to parts of the Dirichlet boundary D . To

this end, we define by F the shrinking set $F := \{z > 0\}$ until the end of this work and obtain the following result.

Lemma 3.7 (cf. [HK12, Lemma 4.18]) *There exists a function $u \in L_t^2 H_{x,\text{loc}}^1(\mathfrak{A}(F); \mathbb{R}^n)$ such that $\epsilon(u) = \widehat{e}$ a.e. in $\mathfrak{A}_\Gamma(F)$ and $u = b$ on the boundary $D_T \cap \mathfrak{A}_\Gamma(F)$.*

A related result can be shown for the sequence $\{\mu_\varepsilon\}$ by exploiting the estimates in Lemma 3.6. To proceed, we recall a definition introduced in [HK12].

Definition 3.8 (cf. [HK12, Definition 4.1]) *Let $H \subseteq \overline{\Omega}$ be a relatively open subset. We call a countable family $\{U_k\}$ of open sets $U_k \subset\subset H$ a fine representation for H if for every $x \in H$ there exist an open set $U \subseteq \mathbb{R}^n$ with $x \in U$ and an $k \in \mathbb{N}$ such that $U \cap \Omega \subseteq U_k$.*

Lemma 3.9 *Let a sequence $\{t_m\} \subseteq [0, T]$ containing T be dense. There exist a fine representation $\{U_k^m\}_{k \in \mathbb{N}}$ for $F(t_m)$ for every $m \in \mathbb{N}$, a function $\mu \in L_t^2 H_{x,\text{loc}}^1(F)$ and a subsequence of $\{\mu_\varepsilon\}$ (also denoted by $\{\mu_\varepsilon\}$) such that for all $k, m \in \mathbb{N}$*

$$\mu_\varepsilon \rightharpoonup \mu \text{ in } L^2(0, t_m; H^1(U_k^m)) \quad (33)$$

as $\varepsilon \rightarrow 0^+$.

Proof. A fine representation $\{U_k^m\}_{k \in \mathbb{N}}$ of $F(t_m)$ can be constructed by countably many open cubes $Q \subset\subset F(t_m) \cap \Omega$ together with finitely many open sets of the form $U \cap \Omega$ such that U satisfies (27) from Lemma 3.6 (ii). Then, we have for each $k, m \in \mathbb{N}$ the estimate

$$\|\mu_\varepsilon\|_{L^2(0, t; H^1(U_k^m))} \leq C$$

for all $0 < \varepsilon \ll 1$. By successively choosing sub-sequences and by a diagonal argument, we obtain a $\mu \in L_t^2 H_{x,\text{loc}}^1(F)$ such that (33) is satisfied (cf. proof of [HK12, Lemma 4.18]). \square

The a-priori estimates and the convergence properties of $\{z_\varepsilon\}$ in Lemma 3.3 and of $\{\mu_\varepsilon\}$ in Lemma 3.9, respectively, yield the following corollary.

Corollary 3.10 *It holds for $\varepsilon \rightarrow 0^+$:*

$$\begin{aligned} m(z_\varepsilon) \nabla \mu_\varepsilon &\rightharpoonup m(z) \nabla \mu \text{ in } L^2(F; \mathbb{R}^n), \\ m(z_\varepsilon) \nabla \mu_\varepsilon &\rightarrow 0 \text{ in } L^2(\Omega_T \setminus F; \mathbb{R}^n). \end{aligned}$$

Now, we have all necessary convergence properties to perform the degenerate limit in (10)-(15). To proceed, we need the following auxiliary result.

Lemma 3.11 *Let $\{t_m\}$ and $\{U_k^m\}$ be as in Lemma 3.9. Then, for every compact subset $K \subseteq F$ there exist a finite set $I \subseteq \mathbb{N}$, values $m_k \in \mathbb{N}$, $k \in I$, and functions $\psi_k \in C^\infty(\overline{\Omega_T})$, $k \in I$, such that*

- (i) $K \cap \Omega_T \subseteq \bigcup_{k \in I} U_k^{m_k} \times (0, t_{m_k})$,
- (ii) $\text{supp}(\psi_k) \subseteq \overline{U_k^{m_k}} \times [0, t_{m_k}]$,
- (iii) $\sum_{k \in I} \psi_k \equiv 1$ on K .

Proof. We extend the family of open sets $\{V_k^m\}$ given by $V_k^m := U_k^m \times (0, t_{m_k})$ in the following way. Define

$$\mathcal{P} := \left\{ \{W_k^m\}_{k,m \in \mathbb{N}} \mid W_k^m \subseteq \mathbb{R}^{n+1} \text{ is open with } W_k^m \cap \Omega_T = U_k^m \times (0, t_{m_k}) \right\}.$$

We see that \mathcal{P} is non-empty and that every totally ordered subset of \mathcal{P} has an upper bound with respect to the " \leq " ordering defined by

$$\{W_k^m\} \leq \{\widetilde{W}_k^m\} \Leftrightarrow W_k^m \subseteq \widetilde{W}_k^m \text{ for all } k, m \in \mathbb{N}.$$

By Zorn's lemma, we find a maximal element $\{\widetilde{V}_k^m\}$. It holds

$$F \subseteq \bigcup_{k,m \in \mathbb{N}} \widetilde{V}_k^m. \quad (34)$$

Assume that this condition fails. Then, because of $F \cap \Omega_T = \bigcup_{k,m \in \mathbb{N}} V_k^m$, there exists a $p = (x, t) \in F \cap \partial(\Omega_T)$ with $p \notin \bigcup_{k,m \in \mathbb{N}} \widetilde{V}_k^m$.

Let us consider the case $t < T$. Since $F \subseteq \overline{\Omega_T}$ is relatively open, we find an $m_0 \in \mathbb{N}$ with $x \in F(t_{m_0})$ and $t_{m_0} > t$. By the fine representation property of $\{U_k^{m_0}\}_{k \in \mathbb{N}}$ for $F(t_{m_0})$, we find an open set $U \subseteq \mathbb{R}^n$ with $x \in U$ and $k_0 \in \mathbb{N}$ such that $U \cap \Omega \subseteq U_{k_0}^{m_0}$.

The family $\{\widetilde{W}_k^m\}$ given by

$$\widetilde{W}_k^m := \begin{cases} \widetilde{V}_k^m \cup U \times (-\infty, t_{m_0}) & \text{if } k = k_0 \text{ and } m = m_0, \\ \widetilde{V}_k^m & \text{else,} \end{cases}$$

satisfies $\{\widetilde{W}_k^m\} \in \mathcal{P}$ and $p \in \bigcup_{k,m \in \mathbb{N}} \widetilde{W}_k^m$ which contradicts the maximality property of $\{\widetilde{V}_k^m\}$.

In the case $t = T$, we also find $k_0, m_0 \in \mathbb{N}$ and an open set $U \subseteq \mathbb{R}^n$ with $x \in U$ such that $U \cap \Omega \subseteq U_{k_0}^{m_0}$ and $t_{m_0} = T$. The family $\{\widetilde{W}_k^m\}$ given by

$$\widetilde{W}_k^m := \begin{cases} \widetilde{V}_k^m \cup U \times \mathbb{R} & \text{if } k = k_0 \text{ and } m = m_0, \\ \widetilde{V}_k^m & \text{else,} \end{cases}$$

also contradicts the maximality of $\{\widetilde{V}_k^m\}$. Therefore, (34) is proven.

Heine-Borel theorem yields

$$K \subseteq \bigcup_{k \in I} \widetilde{V}_k^{m_k}$$

for a finite set $I \subseteq \mathbb{N}$ and values $m_k \in \mathbb{N}$, $k \in I$. Together with a partition of unity argument, we get functions $\psi_k \in C^\infty(\overline{\Omega_T})$ such that (i)-(iii) hold. \square

The degenerate limit $\varepsilon \rightarrow 0^+$ can be performed as follows:

- We define the strain by $e := \widehat{e}|_F \in L^2(F; \mathbb{R}^{n \times n})$ and obtain for the remaining variables

$$\begin{aligned} c &\in L^\infty(0, T; H^1(\Omega)) \cap H^1(0, T; (H^1(\Omega))^*), & u &\in L_t^2 H_{x,\text{loc}}^1(\mathfrak{A}_D(F); \mathbb{R}^n), \\ z &\in L^\infty(0, T; W^{1,p}(\Omega)) \cap H^1(0, T; L^2(\Omega)), & \mu &\in L_t^2 H_{x,\text{loc}}^1(F) \end{aligned}$$

with $e = \epsilon(u)$ in $\mathfrak{A}(F)$.

- Passing to the limit $\varepsilon \rightarrow 0^+$ in (10), (13) and (15) imply properties (5), (8) and (9), cf. [HK12].
- Using Lemma 3.3 (a) and Corollary 3.10, we can pass to $\varepsilon \rightarrow 0^+$ in (11) and obtain (6).
- Let $\zeta \in L^2(0, T; H^1(\Omega))$ with $\text{supp}(\zeta) \subseteq F$ be a test-function. Furthermore, let $\{\psi_l\}$ be a partition of unity of the compact set $K := \text{supp}(\zeta)$ according to Lemma 3.11. For each $l \in \mathbb{N}$, we obtain $\text{supp}(\zeta\psi_l) \subseteq \overline{U_l^{m_l}} \times [0, t_{m_l}]$. Then, integrating (12) in time from 0 to t_{m_l} , testing the result with $\zeta\psi_l$ and passing to $\varepsilon \rightarrow 0^+$ by using Lemma 3.3 and Lemma 3.9 show

$$\int_0^{t_m} \int_{\Omega} \mu \zeta \psi_l \, dx \, ds = \int_0^{t_m} \int_{\Omega} \left(\nabla c \cdot \nabla(\zeta \psi_l) + \Psi_{,c}(c) \zeta \psi_l + W_{,c}(c, \widehat{e}, z) \zeta \psi_l \right) \, dx \, ds.$$

Summing with respect to $l \in I$ and noticing $\sum_{l \in I} \psi_l \equiv 1$ on $\text{supp}(\zeta)$ yield (7).

In conclusion, the limit procedure in this section yields functions (c, e, u, z, μ) with $e = \epsilon(u)$ in $\mathfrak{A}_D(F)$ which satisfy properties (ii)-(v) of Definition 2.1. In particular, the damage function z has no jumps with respect to time. However, we cannot ensure that $\{z > 0\}$ equals $\mathfrak{A}_D(\{z > 0\})$ and, moreover, if $F \setminus \mathfrak{A}_D(\{z > 0\}) \neq \emptyset$, it is not clear whether u can be extended to a function on F such that $e = \epsilon(u)$ also holds in F . This issue is addressed in the next section where such limit functions are concatenated in order to obtain global-in-time weak solutions with fineness $\eta > 0$ by Zorn's lemma.

4 Existence results

In this section, we are going to prove the main results of this paper.

Theorem 4.1 (Maximal local-in-time weak solutions)

Let $b \in W^{1,1}(0, T; W^{1,\infty}(\Omega; \mathbb{R}^n))$, $c^0 \in H^1(\Omega)$ and $z^0 \in W^{1,p}(\Omega)$ with $0 < \kappa \leq z^0 \leq 1$ in Ω be initial-boundary data. Then there exist a maximal value $\widehat{T} > 0$ with $\widehat{T} \leq T$ and functions c, u, z, μ defined on the time interval $[0, \widehat{T}]$ such that (c, u, z, μ) is a weak solution according to Definition 2.1. Therefore, if $\widehat{T} < T$, (c, u, z, μ) cannot be extended to a weak solution on $[0, \widehat{T} + \varepsilon]$.

Proof. Zorn's lemma can be applied to the set

$$\mathcal{P} := \{(\widehat{T}, c, u, z, \mu) \mid 0 < \widehat{T} \leq T \text{ and } (c, u, z, \mu) \text{ is a weak solution on } [0, \widehat{T}] \text{ according to Definition 2.1}\}$$

to find a maximal element with respect to the following partial ordering

$$\begin{aligned} (\widehat{T}_1, c_1, u_1, z_1, \mu_1) \leq (\widehat{T}_2, c_2, u_2, z_2, \mu_2) &\Leftrightarrow \widehat{T}_1 \leq \widehat{T}_2, \, c_2|_{[0, \widehat{T}_1]} = c_1, \, u_2|_{[0, \widehat{T}_1]} = u_1, \\ &z_2|_{[0, \widehat{T}_1]} = z_1, \, \mu_2|_{[0, \widehat{T}_1]} = \mu_1. \end{aligned} \tag{35}$$

Indeed, $\mathcal{P} \neq \emptyset$ by the result in Section 3. More precisely, since $z \in L^\infty(0, T; W^{1,p}(\Omega)) \cap H^1(0, T; L^2(\Omega))$ and since $0 < \kappa \leq z^0$, we find an $\varepsilon > 0$ such that $\{z(t) > 0\} = \mathfrak{A}_D(\{z(t) > 0\})$ for all $t \in [0, \varepsilon]$. For the proof that every totally ordered subset of \mathcal{P} has an upper bound, we refer to [HK12]. \square

The proof of global-in-time existence of weak solutions with fineness η requires a concatenation property (see Lemma 4.4) which is, in turn, based on some deeper insights into the Γ -limit \mathfrak{E} introduced in Section 3. To this end, it is necessary to have more information about the recovery sequences for $\mathfrak{F}_\varepsilon \xrightarrow{\Gamma} \mathfrak{F}$.

We will introduce the following substitution method: Assume that $u \in H^1(\Omega; \mathbb{R}^n)$ minimizes $\mathcal{F}_\varepsilon(c, \epsilon(\cdot), z)$ with Dirichlet data ξ on D . Then, by expressing the elastic energy density W in terms of its derivative $W_{,e}$, i.e.

$$W(c, e, z) = \frac{1}{2} W_{,e}(c, e, z) : e + \frac{1}{2} z \varphi^2(c) : e + z \varphi^3(c),$$

and by testing the momentum balance equation with $\zeta = u - \tilde{u}$ for a function $\tilde{u} \in H^1(\Omega; \mathbb{R}^n)$ with $\tilde{u} = \xi$ on D , the elastic energy term in \mathcal{F}_ε can be rewritten as

$$\int_{\Omega} W_\varepsilon(c, \epsilon(u), z) \, dx = \int_{\Omega} (g(z) + \varepsilon) \left(\varphi^1 \epsilon(u) : \epsilon(\tilde{u}) + \frac{1}{2} \varphi^2(c) : (\epsilon(u) + \epsilon(\tilde{u})) + \varphi^3(c) \right) \, dx. \quad (36)$$

For convenience, in the following proof, we define the density function \widetilde{W} as

$$\widetilde{W}_\varepsilon(c, e, e_1, z) := (g(z) + \varepsilon) \left(\varphi^1 e : e_1 + \frac{1}{2} \varphi^2(c) : (e + e_1) + \varphi^3(c) \right).$$

Lemma 4.2 *For every $c \in H^1(\Omega)$, $\xi \in W^{1,\infty}(\Omega)$ and $z \in W^{1,p}(\Omega)$ there exists a sequence $\delta_\varepsilon \rightarrow 0^+$ such that $(c, \xi, (z - \delta_\varepsilon)^+) \rightarrow (c, \xi, z)$ is a recovery sequence for $\mathfrak{F}_\varepsilon \xrightarrow{\Gamma} \mathfrak{F}$.*

Proof. We follow the idea of the proof in [HK12, Lemma 4.9]. But here we have to deal with the additional concentration variable which complicates the calculation. Let $(c_\varepsilon, \xi_\varepsilon, z_\varepsilon) \rightarrow (c, \xi, z)$ be a recovery sequence and $\delta_\varepsilon \rightarrow 0^+$ such that $(z - \delta_\varepsilon)^+ \leq z_\varepsilon$. Consider

$$\mathfrak{F}_\varepsilon(c, \xi, (z - \delta_\varepsilon)^+) - \mathfrak{F}_\varepsilon(c_\varepsilon, \xi_\varepsilon, z_\varepsilon) = \underbrace{\mathfrak{F}_\varepsilon(c, \xi, (z - \delta_\varepsilon)^+) - \mathfrak{F}_\varepsilon(c, \xi, z_\varepsilon)}_{A_\varepsilon} + \underbrace{\mathfrak{F}_\varepsilon(c, \xi, z_\varepsilon) - \mathfrak{F}_\varepsilon(c_\varepsilon, \xi_\varepsilon, z_\varepsilon)}_{B_\varepsilon}.$$

We have $A_\varepsilon \leq 0$ since $\mathcal{F}_\varepsilon(c, \epsilon(\xi + \zeta), (z - \delta_\varepsilon)^+) \leq \mathcal{F}_\varepsilon(c, \epsilon(\xi + \zeta), z_\varepsilon)$ for all $\zeta \in H_D^1(\Omega; \mathbb{R}^n)$. Now, we focus on the second term of the right hand side. Let $u_\varepsilon, v_\varepsilon \in H_D^1(\Omega; \mathbb{R}^n)$ be given by

$$u_\varepsilon = \arg \min_{\zeta \in H_D^1(\Omega; \mathbb{R}^n)} \mathcal{F}_\varepsilon(c, \epsilon(\xi + \zeta), z_\varepsilon), \quad v_\varepsilon = \arg \min_{\zeta \in H_D^1(\Omega; \mathbb{R}^n)} \mathcal{F}_\varepsilon(c_\varepsilon, \epsilon(\xi_\varepsilon + \zeta), z_\varepsilon).$$

By using the substitution (36) for $W_\varepsilon(c, \epsilon(\xi + u_\varepsilon), z_\varepsilon)$ with test-function $\tilde{u} = v_\varepsilon$ and (36) for $W_\varepsilon(c_\varepsilon, \epsilon(\xi_\varepsilon + v_\varepsilon), z_\varepsilon)$ with test-function $\tilde{u} = u_\varepsilon$, we obtain a calculation as follows:

$$\begin{aligned} B_\varepsilon &= \mathcal{F}_\varepsilon(c, \epsilon(\xi + u_\varepsilon), z_\varepsilon) - \mathcal{F}_\varepsilon(c_\varepsilon, \epsilon(\xi_\varepsilon + v_\varepsilon), z_\varepsilon) \\ &= \int_{\Omega} \left(\widetilde{W}(c, \epsilon(\xi + u_\varepsilon), \epsilon(\xi + v_\varepsilon), z_\varepsilon + \varepsilon) - \widetilde{W}(c_\varepsilon, \epsilon(\xi_\varepsilon + v_\varepsilon), \epsilon(\xi_\varepsilon + u_\varepsilon), z_\varepsilon + \varepsilon) \right) \, dx \\ &= \int_{\Omega} (g(z_\varepsilon) + \varepsilon) \left(\varphi^1 \epsilon(\xi + u_\varepsilon) : \epsilon(\xi + v_\varepsilon) - \varphi^1 \epsilon(\xi_\varepsilon + u_\varepsilon) : \epsilon(\xi_\varepsilon + v_\varepsilon) \right. \\ &\quad \left. + \frac{1}{2} \varphi^2(c) : \epsilon(2\xi + u_\varepsilon + v_\varepsilon) - \frac{1}{2} \varphi^2(c_\varepsilon) : \epsilon(2\xi + v_\varepsilon + u_\varepsilon) + \varphi^3(c) - \varphi^3(c_\varepsilon) \right) \, dx \end{aligned}$$

$$\begin{aligned}
&= \int_{\Omega} (g(z_{\varepsilon}) + \varepsilon) \left(\varphi^1 \epsilon(\xi) : \epsilon(\xi) - \varphi^1 \epsilon(\xi_{\varepsilon}) : \epsilon(\xi_{\varepsilon}) + \varphi^1 \epsilon(u_{\varepsilon} + v_{\varepsilon}) : \epsilon(\xi - \xi_{\varepsilon}) \right. \\
&\quad \left. + \varphi^2(c) : \epsilon(\xi - \xi_{\varepsilon}) + \frac{1}{2}(\varphi^2(c) - \varphi^2(c_{\varepsilon})) : \epsilon(2\xi_{\varepsilon} + u_{\varepsilon} + v_{\varepsilon}) + \varphi^3(c) - \varphi^3(c_{\varepsilon}) \right) dx \\
&\leq \int_{\Omega} (g(z_{\varepsilon}) + \varepsilon) \left(\varphi^1 \epsilon(\xi) : \epsilon(\xi) - \varphi^1 \epsilon(\xi_{\varepsilon}) : \epsilon(\xi_{\varepsilon}) + \varphi^2(c) : \epsilon(\xi - \xi_{\varepsilon}) + \varphi^3(c) - \varphi^3(c_{\varepsilon}) \right) dx \\
&\quad + \|(g(z_{\varepsilon}) + \varepsilon) \varphi^1 \epsilon(u_{\varepsilon} + v_{\varepsilon})\|_{L^2(\Omega)} \|\epsilon(\xi - \xi_{\varepsilon})\|_{L^2(\Omega)} \\
&\quad + \frac{1}{2} \|\varphi^2(c) - \varphi^2(c_{\varepsilon})\|_{L^2(\Omega)} \left(\|(g(z_{\varepsilon}) + \varepsilon) \epsilon(\xi_{\varepsilon} + u_{\varepsilon})\|_{L^2(\Omega)} + \|(g(z_{\varepsilon}) + \varepsilon) \epsilon(\xi_{\varepsilon} + v_{\varepsilon})\|_{L^2(\Omega)} \right)
\end{aligned}$$

Using the convergence properties $c_{\varepsilon} \rightharpoonup c$ in $H^1(\Omega)$, $\xi_{\varepsilon} \rightarrow \xi$ in $W^{1,\infty}(\Omega)$, $z_{\varepsilon} \rightharpoonup z$ in $W^{1,p}(\Omega)$ and the boundedness of $\mathcal{F}_{\varepsilon}(c, \epsilon(\xi + u_{\varepsilon}), z_{\varepsilon})$ and $\mathcal{F}_{\varepsilon}(c_{\varepsilon}, \epsilon(\xi_{\varepsilon} + v_{\varepsilon}), z_{\varepsilon})$ with respect to ε , we conclude $\limsup_{\varepsilon \rightarrow 0^+} B_{\varepsilon} \leq 0$. Now we can proceed as in the proof of [HK12, Lemma 4.9] and the claim follows. \square

Remark 4.3 The knowledge of such recovery sequences for $\mathfrak{F}_{\varepsilon} \xrightarrow{\Gamma} \mathfrak{F}$ gives also more information about \mathfrak{E} . In particular, we obtain that the result in Lemma 4.2 with $\mathfrak{E}_{\varepsilon} \xrightarrow{\Gamma} \mathfrak{E}$ instead of $\mathfrak{F}_{\varepsilon} \xrightarrow{\Gamma} \mathfrak{F}$ also holds true and, moreover, that the following properties are satisfied (cf. [HK12, Corollary 4.10, Lemma 4.11]):

- $\mathfrak{E}(c, \xi, \mathbb{1}_F z) \leq \mathfrak{E}(c, \xi, z)$ $\forall c \in H^1(\Omega), \forall \xi \in W^{1,\infty}(\Omega; \mathbb{R}^n), \forall z \in W^{1,p}(\Omega)$
 $\forall F \subseteq \Omega$ open with $\mathbb{1}_F z \in W^{1,p}(\Omega)$,
- $\mathfrak{E}(c, \xi, z) \leq \mathcal{E}(c, \epsilon(u), z)$ $\forall c \in H^1(\Omega), \forall \xi \in W^{1,\infty}(\Omega; \mathbb{R}^n), \forall z \in W^{1,p}(\Omega)$ with
 $0 \leq z \leq 1, \forall u \in H_{\text{loc}}^1(\{z > 0\}; \mathbb{R}^n)$ with $u = \xi$ on $D \cap \{z > 0\}$.

Lemma 4.4 Let $t_1 < t_2 < t_3$ be real numbers and let $\eta > 0$. Suppose that

- $\tilde{q} := (\tilde{c}, \tilde{e}, \tilde{u}, \tilde{z}, \tilde{\mu}, \tilde{F})$ is a weak solution with fineness η on $[t_1, t_2]$,
- $\hat{q} := (\hat{c}, \hat{e}, \hat{u}, \hat{z}, \hat{\mu}, \hat{F})$ is a weak solution with fineness η on $[t_2, t_3]$ with $\hat{\mathbf{c}}_{t_2}^+ = \mathfrak{E}(\hat{c}(t_2), \hat{b}(t_2), \hat{z}^+(t_2))$
(the value $\mathfrak{c}_{t_2}^+$ for the weak solution \hat{q} in Definition 2.1).

Furthermore, suppose the compatibility condition $\hat{c}(t_2) = \tilde{c}(t_2)$ and $\hat{z}^+(t_2) = \tilde{z}^-(t_2) \mathbb{1}_{\mathfrak{A}_D(\{\tilde{z}^-(t_2) > 0\})}$ and the Dirichlet boundary data $b \in W^{1,1}(t_1, t_3; W^{1,\infty}(\Omega; \mathbb{R}^n))$.

Then, we obtain that $q := (c, e, u, z, \mu, F)$ defined as $q|_{[t_1, t_2]} := \tilde{q}$ and $q|_{[t_2, t_3]} := \hat{q}$ is a weak solution with fineness η on $[t_1, t_3]$.

Proof. Because of the properties in Remark 4.3 we can prove the following crucial energy estimate at time point t_2 :

$$\begin{aligned}
\lim_{s \rightarrow t_2^-} \operatorname{ess\,inf}_{\tau \in (s, t_2)} \mathcal{E}(c(\tau), e(\tau), z(\tau)) &= \lim_{s \rightarrow t_2^-} \operatorname{ess\,inf}_{\tau \in (s, t_2)} \mathcal{E}(c(\tau), e(\tau), z^-(\tau)) \\
&\geq \lim_{s \rightarrow t_2^-} \operatorname{ess\,inf}_{\tau \in (s, t_2)} \mathcal{E}(c(\tau), \epsilon(u(\tau)), z^-(\tau) \mathbb{1}_{\mathfrak{A}_D(\{z^-(\tau) > 0\})}) \\
&\geq \lim_{s \rightarrow t_2^-} \operatorname{ess\,inf}_{\tau \in (s, t_2)} \mathfrak{E}(c(\tau), b(\tau), z^-(\tau) \mathbb{1}_{\mathfrak{A}_D(\{z^-(\tau) > 0\})})
\end{aligned}$$

$$\begin{aligned}
&\geq \mathfrak{E}(c(t_2), b(t_2), \chi) \\
&\geq \mathfrak{E}(c(t_2), b(t_2), z^+(t_2))
\end{aligned}$$

with $\chi := z^-(t_2)\mathbb{1}_{\cap_{\tau \in (t_1, t_2)} \mathfrak{A}_D(\{z^-(\tau) > 0\})}$. With this estimate, we can verify the claim by the same argumentation as for [HK12, Lemma 4.21]. \square

Theorem 4.5 (Global-in-time weak solutions with fineness η)

Let $b \in W^{1,1}(0, T; W^{1,\infty}(\Omega; \mathbb{R}^n))$, $c^0 \in H^1(\Omega)$ and $z^0 \in W^{1,p}(\Omega)$ with $0 \leq z^0 \leq 1$ in Ω and $\{z^0 > 0\}$ admissible with respect to D be initial-boundary data. Furthermore, let $\eta > 0$. Then there exists a weak solution (c, e, u, z, μ) with fineness $\eta > 0$ according to Definition 2.3.

Proof. This result can also be proven by using Zorn's lemma on the set

$$\mathcal{P} := \{(\widehat{T}, c, e, u, z, \mu, F) \mid 0 < \widehat{T} \leq T \text{ and } (c, e, u, z, \mu, F) \text{ is a weak solution with fineness } \eta \text{ on } [0, \widehat{T}] \text{ according to Definition 2.3}\}$$

with an ordering analogously to (35). The assumptions for Zorn's lemma can be proven as in Theorem 4.1 (see [HK12, Proof of Theorem 4.1]). To show that a maximal element from \mathcal{P} is actually a weak solution with fineness η on the time-interval $[0, T]$, we need the concatenation property in Lemma 4.4. Indeed, if a maximal element \tilde{q} is only defined on a time-interval $[0, \tilde{T}]$ with $\tilde{T} < T$ we can apply the degenerated limit procedure in Section 3 to the initial values $c(\tilde{T})$ and $z(\tilde{T})$ to obtain a new limit function \hat{q} . By exploiting Lemma 4.4, \hat{q} is a weak solution with fineness η on the time-interval $[0, \tilde{T} + \varepsilon]$ for a small $\varepsilon > 0$ which contradicts the maximality of \tilde{q} . \square

References

- [BB99] J.W. Barrett and J.F. Blowey. Finite element approximation of the Cahn-Hilliard equation with concentration dependent mobility. *Math. of Computation*, 68(226):487–517, 1999.
- [BCD⁺02] E. Bonetti, P. Colli, W. Dreyer, G. Gilardi, G. Schimperna, and J. Sprekels. On a model for phase separation in binary alloys driven by mechanical effects. *Physica D*, 165:48–65, 2002.
- [BK98] S. M. Buckley and P. Koskela. New Poincaré inequalities from old. *Annales Academiæ Scientiarum Fennicæ Mathematica*, 23:251–260, 1998.
- [BM10] S. Bartels and R. Müller. A posteriori error controlled local resolution of evolving interfaces for generalized Cahn–Hilliard equations. *Interfaces and Free Boundaries*, 12(1):45–73, 2010.
- [BMR09] G. Bouchitte, A. Mielke, and T. Roubíček. A complete-damage problem at small strains. *ZAMP Z. Angew. Math. Phys.*, 60:205–236, 2009.
- [BP05] L. Bartkowiak and I. Pawłow. The Cahn-Hilliard-Gurtin system coupled with elasticity. *Control and Cybernetics*, 34:1005–1043, 2005.
- [BS04] E. Bonetti and G. Schimperna. Local existence for Frémond's model of damage in elastic materials. *Contin. Mech. Thermodyn.*, 16(4):319–335, 2004.
- [CH58] J.W. Cahn and J.E. Hilliard. Free energy of a uniform system. i. interfacial free energy. *Journal of Chemical Physics*, 28:258–267, 1958.
- [CMP00] M. Carrive, A. Miranville, and A. Piétrus. The Cahn-Hilliard equation for deformable elastic continua. *Adv. Math. Sci. Appl.*, 10(2):539–569, 2000.
- [DM01] W. Dreyer and W.H. Müller. Modeling diffusional coarsening in eutectic tinlead solders: A quantitative approach. *Internat. J. Solids and Structures*, 38(8):1433–1458, 2001.

- [EG96] C.M. Elliott and H. Garcke. On the Cahn-Hilliard equation with degenerate mobility. *SIAM J. Math. Anal.*, 27:404–423, 1996.
- [Ell89] C.M. Elliott. The Cahn-Hilliard Model for the Kinetics of Phase Separation. In *Mathematical Models for Phase Change Problems* (ed. J.-F. Rodrigues), volume 88 of *International Series of Numerical Mathematics*, pages 35–73. Birkhäuser Basel, 1989.
- [FK09] J. R. Fernández and K. L. Kuttler. An existence and uniqueness result for an elasto-piezoelectric problem with damage. *Math. Mod. Meth. Appl. Sci.*, 19(1):31–50, 2009.
- [Gar00] H. Garcke. *On mathematical models for phase separation in elastically stressed solids*. Habilitation thesis, University Bonn, 2000.
- [Gar05a] H. Garcke. Mechanical Effects in the Cahn-Hilliard Model: A Review on Mathematical Results. In A. Miranville, editor, *Mathematical Methods and Models in phase transitions*, pages 43–77. Nova Science Publ., 2005.
- [Gar05b] H. Garcke. On a Cahn-Hilliard model for phase separation with elastic misfit. *Annales de l’Institut Henri Poincaré (C) Non Linear Analysis*, 22(2):165 – 185, 2005.
- [Gia05] A. Giacomini. Ambrosio–Tortorelli approximation of quasi-static evolution of brittle fractures. *Calc. Var. Partial Differ. Equ.*, 22(2):129–172, 2005.
- [GRW01] H. Garcke, M. Rumpf, and U. Weikard. The Cahn-Hilliard equation with elasticity: Finite element approximation and qualitative studies. *Interfaces Free Bound.*, 3:101–118, 2001.
- [GUE⁺07] M.G.D. Geers, R.L.J.M. Ubachs, M. Erinc, M.A. Matin, P.J.G. Schreurs, and W.P. Vellinga. Multiscale Analysis of Microstructura Evolution and Degradation in Solder Alloys. *Internatinal Journal for Multiscale Computational Engineering*, 5(2):93–103, 2007.
- [Gur89] M.E. Gurtin. On a Nonequilibrium Thermodynamics of Capillarity and Phase. *Quarterly of Applied Mathematics*, 47(1):129–145, 1989.
- [HCW91] P. G. Harris, K. S Chaggar, and M. A. Whitmore. The Effect of Ageing on the Microstructure of 60:40 Tin–lead Solders. *Soldering & Surface Mount Technology Improved physical understanding of intermittent failure in continuous*, 3:20–33, 1991.
- [HK11] C. Heinemann and C. Kraus. Existence of weak solutions for Cahn-Hilliard systems coupled with elasticity and damage. *Adv. Math. Sci. Appl.*, 21(2):321–359, 2011.
- [HK12] C. Heinemann and C. Kraus. Complete damage in linear elastic materials — modeling, weak formulation and existence results. *WIAS preprint no. 1722 (submitted)*, 2012.
- [HK13] C. Heinemann and C. Kraus. Existence results for diffuse interface models describing phase separation and damage. *Eur. J. Appl. Math.*, 24(2):179–211, 2013.
- [KRZ11] D. Knees, R. Rossi, and C. Zanini. *A vanishing viscosity approach to a rate-independent damage model*. WIAS preprint no. 1633. WIAS, 2011.
- [Mer05] T. Merkle. The Cahn-Larché system: A model for spinodal decomposition in eutectic solder; modelling, analysis and simulation. PhD-thesis, Universität Stuttgart, Stuttgart, 2005.
- [Mie11] A. Mielke. Complete-damage evolution based on energies and stresses. *Discrete Contin. Dyn. Syst., Ser. S*, 4(2):423–439, 2011.
- [MR06] A. Mielke and T. Roubíček. Rate-independent damage processes in nonlinear elasticity. *Mathematical Models and Methods in Applied Sciences*, 16:177–209, 2006.
- [MRZ10] A. Mielke, T. Roubíček, and J. Zeman. Complete Damage in elastic and viscoelastic media. *Comput. Methods Appl. Mech. Engrg*, 199:1242–1253, 2010.
- [MT10] A. Mielke and M. Thomas. Damage of nonlinearly elastic materials at small strain — Existence and regularity results. *ZAMM Z. Angew. Math. Mech*, 90:88–112, 2010.
- [PZ08] I. Pawłó and W. M. Zajączkowski. Measure-valued solutions of a heterogeneous Cahn-Hilliard system in elastic solids. *Colloquium Mathematicum*, 112 No.2, 2008.
- [RR12] E. Rocca and R. Rossi. A degenerating PDE system for phase transitions and damage. *arXiv:1205.3578v1*, 2012.
- [Sim86] J. Simon. Compact sets in the space $L^p(0, T; B)$. *Annali di Matematica Pura ed Applicata*, 146:65–96, 1986.
- [USG07] R.L.J.J. Ubachs, P.J.G. Schreurs, and M.G.D. Geers. Elasto-viscoplastic nonlocal damage modelling of thermal fatigue in anisotropic lead-free solder. *Mechanics of Materials*, 39:685–701, 2007.
- [Wei01] U. Weikard. Numerische Lösungen der Cahn-Hilliard-Gleichung und der Cahn-Larché-Gleichung. PhD-thesis, Universität Bonn, Bonn, 2001.
- [Zie89] W.P. Ziemer. *Weakly differentiable functions*. Springer-Verlag New York, Inc., 1989.